

**Decay Estimates For the Solutions of Linear and Nonlinear  
Damped Wave Equation**

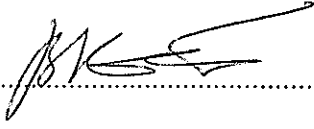
**by**

**Neslihan Nesliye AYKIR**

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Prof. Varga K. Kalantarov.....

Prof. Ali Ulger.....

Assoc. Prof. Mustafa Polat.....

Koc University  
October 2, 2011

## **Abstract**

In the thesis, we investigate the asymptotic behavior of solutions of linear and nonlinear wave equations. In the first chapter, we give some definitions, lemmas and theorems that will be used. In the second chapter, we establish the decay estimates for solutions of a linear wave equation with mixed boundary conditions. In the third chapter, we study the asymptotic behavior of solutions of initial boundary value problem for the wave equation with nonlinear damping term. In the fourth chapter we derive the decay estimate for solution of the initial boundary value problem for the equation of nonlinear vibrations of an elastic string. What is shown by energy decay for all of these equations is the stability of these equations. As a result, we show the stability of those three different wave equations.

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# Chapter 1

## Introduction

This thesis is devoted to the study of asymptotic behavior of solutions to initial boundary value problems for damped linear and nonlinear wave equations.

The first problem we consider is the following:

$$u_{tt} - \Delta u + qu = 0, \quad x \in \Omega, t \in \mathbb{R}^+$$

$$u(x, t) = 0, \quad x \in \Gamma_0, t \in \mathbb{R}^+$$

$$\partial_\nu u + au + lu_t = 0, \quad x \in \Gamma_1, t \in \mathbb{R}^+$$

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad x \in \Omega.$$

In the second chapter, following Komornik <sup>[Kom1]</sup><sub>[3]</sub> decay estimate for classical solution of this problem is derived.

In the third chapter by using the so called Nakao inequality we investigate asymptotic behavior of solutions to the following problem

$$\begin{cases} u_{tt} - \sum_{i,j=1}^n (a_{ij}(x)u_{x_j})_{x_i} + a_0(x)u + \beta(x, u_t) = f(x, t), & x \in \Omega, t > 0, \\ u(x, t) = 0, & x \in \partial\Omega, t > 0, \\ u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), & x \in \Omega, \end{cases}$$

where  $\Omega$  is a bounded domain with a smooth boundary  $\partial\Omega$ ,

1.  $a_{ij}(x)$  and  $a_0(x)$  are measurable and bounded functions on  $\Omega$  :

$$a_{ij}(x) = a_{ji}(x), \quad a_0(x) \geq 0 \quad \forall x \in \Omega$$

and

$$\sum_{i,j=1}^n a_{ij}(x)\xi_i\xi_j \geq \nu^{-1}|\xi|^2, \quad \forall \xi \in \mathbb{R}^n$$

where  $\nu > 0$

2.  $\beta(x, z)$  is measurable on  $\Omega \times \mathbb{R}$

$$\beta(x, 0) = 0, \quad |\beta(x, z_1) - \beta(x, z_2)| \leq k_0(1 + |z_1|^\gamma + |z_2|^\gamma)|z_1 - z_2|,$$

$$(\beta(x, z_1) - \beta(x, z_2))(z_1 - z_2) \geq k_1|z_1 - z_2|^{\gamma+2}$$

3.  $0 \leq \gamma \leq \frac{4}{n-2}$  if  $n \geq 3$   $0 \leq \gamma \leq \infty$  if  $n = 1, 2$ .

Following Nakao [5], the difference between bounded solution and any solution of this problem tend to zero with a polynomial rate as  $t \rightarrow \infty$  is shown.

The bounded and the almost periodic solutions for this problem was investigated by Amerio and Prouse [15]. They proved that for any solution  $v(t)$  the bounded solution or almost periodic solution  $u(t)$  satisfies

$$\lim_{t \rightarrow \infty} \|u(t) - v(t)\|_E = 0.$$

Here  $\|\cdot\|_E$  denotes the energy norm. They also proved that for each initial data  $(u_0, u_1) \in H_0^1(\Omega) \times L^2(\Omega)$ , the problem has a unique weak solution and if  $f(x, t)$  is uniformly continuous then the problem has unique bounded weak solution.

In the fourth chapter, following papers of K. Nishihara we study the asymptotic behavior of solutions of the problem:

$$\begin{cases} u_{tt} + (1 + \|A^{\frac{1}{2}}u(t)\|^2)Au(t) + 2\gamma u_t = f(x, t), & x \in \Omega, t > 0, \\ u(x, t) = 0, & x \in \partial\Omega, t > 0, \\ u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), & x \in \Omega, \end{cases}$$

The assumption on  $A$  is self adjoint, positive definite operator with discrete spectrum and  $A^{-1}$  is compact.

To show the asymptotic behavior of that equation, in Nishihara [7]

$$\|F(A^{\frac{1}{2}})u_t\|, \|F(A^{\frac{1}{2}})A^{\frac{1}{2}}u\| \leq C \exp(-\theta t), \quad \text{where } \theta > 0$$

is proven.

Here we also presented the result about continuous dependence of solution to this problem on the data.

The Cauchy and mixed problem for the equation with  $\gamma = 0$  was investigated by a number of authors: Pohozaev [12], Perla [11], Nishida [10], Dickey [9], Rivera [13], Greenberg and Hul [16].

When  $\gamma > 0$  the global existence and asymptotic behavior of the solutions was shown by Yamada [14] if the data are small in some sense. In [7] no smallness condition is assumed on the data, but sufficient regularity is assumed.

## 1.1 Prelimineries

### 1.1.1 Function Spaces

**Definition 1.1.1.** A Banach space is complete linear normed space.

**Definition 1.1.2.** A Hilbert space  $H$  is a complete inner product space.

**Definition 1.1.3.** The function  $v$  is said to be of class  $C^k(G)$  if the derivatives  $v', v'', \dots, v^{(k)}$  exist and are continuous in  $G$ .

**Definition 1.1.4.** The function  $f$  is said to be of class  $C^\infty(G)$ , if it has derivatives of all orders.

**Definition 1.1.5.** Let  $\Omega \in \mathbb{R}^n$  and  $f$  is defined on  $\Omega$  then  $f \in C_0^\infty(\Omega)$  if  $f$  is infinitely differentiable function with compact support.

**Definition 1.1.6.** For  $1 \leq p < \infty$ ,  $L^p(\Omega)$  is the set of all measurable functions  $u(x)$  in  $\Omega$  such that the norm

$$\|u\|_{L^p(\Omega)} \leq \left( \int_{\Omega} |u(x)|^p dx \right)^{\frac{1}{p}}$$

is finite.

**Definition 1.1.7.**  $L^\infty(\Omega)$  is the set of all bounded measurable functions in  $\Omega$ ; the norm is defined by

$$\|u\|_{L^\infty(\Omega)} = \text{ess sup}_{x \in \Omega} |u(x)|.$$

All  $L^p(\Omega)$  spaces are Banach spaces but  $L^2(\Omega)$  is a Hilbert space and the inner product on  $L^2(\Omega)$  is defined as:

$$(u, v) = \left( \int_{\Omega} u(x)v(x)dx \right), \text{ here } u, v \in L^2(\Omega)$$

Assume  $\Omega \subset \mathbb{R}^n$  is bounded,  $u \in C^1(\Omega)$  and  $\Psi \in C_0^\infty(\Omega)$ , by integration by parts we get,

$$\int_{\Omega} u \Psi_{x_i} dx = - \int_{\Omega} u_{x_i} \Psi dx. \quad (1.1) \quad \boxed{\text{wd}}$$

Here the term on  $\partial\Omega$  vanishes since  $\psi$  has compact support.

For the left hand side of the equation (1.1) if we assume that  $u$  is locally integrable on  $\Omega$ , then (1.1) still makes sense. So we can define weak derivative notion as below.

**Definition 1.1.8.** Let  $\alpha$  be multiindex such that  $\alpha = (\alpha_1, \dots, \alpha_n)$ ,  $|\alpha| = \alpha_1 + \dots + \alpha_n$ . Suppose that  $u, v$  be locally integrable function defined on  $\Omega$  and

$$\int_{\Omega} u D^{\alpha} \Psi dx = (-1)^{|\alpha|} \int_{\Omega} v \Psi dx, \forall \Psi \in C_0^{\infty}(\Omega).$$

Then  $v$  is called the weak derivative of  $u$  in  $\Omega$  and denoted by  $D^{\alpha}u$ .

Now we can define the Sobolev space  $H^1(\Omega)$ ,  $H^2(\Omega)$ ,  $H_0^1(\Omega)$  and  $H_{\Gamma_0}^1(\Omega)$ .

**Definition 1.1.9.** The Sobolev space  $H^1(\Omega)$  is the set of all functions from  $L^2(\Omega)$  such that all the first order weak derivatives are also from  $L^2(\Omega)$ . i.e

$$H^1(\Omega) := \{u \in L^2(\Omega) : u_{x_i} \in L^2(\Omega), 1 \leq i \leq n\}.$$

This space is equipped with the norm:

$$\|u\|_{H^1(\Omega)} = \left( \int_{\Omega} \sum_{|\alpha| \leq 1} |D^{\alpha}u|^2 dx \right)^{\frac{1}{2}}.$$

**Definition 1.1.10.** The Sobolev space  $H^2(\Omega)$  is the set of all functions from  $L^2(\Omega)$  such that all the first order weak derivatives and second order weak derivatives are also from  $L^2(\Omega)$ . i.e

$$H^2(\Omega) := \{u \in L^2(\Omega) : u_{x_i}, u_{x_i x_j} \in L^2(\Omega), 1 \leq i, j \leq n\}.$$

This space is equipped with the norm:

$$\|u\|_{H^2(\Omega)} = \left( \int_{\Omega} \sum_{|\alpha| \leq 2} |D^{\alpha}u|^2 dx \right)^{\frac{1}{2}}.$$

**Definition 1.1.11.** The Sobolev space  $H_0^1(\Omega)$  is the closure of the space  $C_0^{\infty}(\Omega)$  with respect to the norm on  $H^1(\Omega)$ . i.e.

$$H_0^1(\Omega) := \{u \in H^1(\Omega) : u_m \rightarrow u \text{ where } (u_m)_{m \in \mathbb{N}} \subset C_0^{\infty}(\Omega)\}.$$

This space is equipped with the norm:

$$\|u\|_{H_0^1(\Omega)} = \left( \int_{\Omega} |\nabla u|^2 dx \right)^{\frac{1}{2}}.$$

**Definition 1.1.12.** Let  $\Gamma_0, \Gamma_1$  be the partition of  $\partial\Omega$ . Then

$$H_{\Gamma_0}^1(\Omega) := \{u \in H^1(\Omega) : u = 0 \text{ on } \Gamma_0 \subset \partial\Omega\}.$$

nak3L **Lemma 1.1.13. (Sobolev Lemma) :** Let  $\gamma$  be ;  $0 \leq \gamma \leq \frac{4}{n-2}$  if  $n \geq 3$   $0 \leq \gamma < \infty$  if  $n = 1, 2$ . Then

$$\|u\|_{L^{\gamma+2}} \leq S_{\gamma+2} \|u\|_{H_0^1(\Omega)} \tag{1.2} \span style="border: 1px solid black; padding: 2px;">nak3$$

for  $u \in H_0^1(\Omega)$  where  $S_{\gamma+2}$  is a constant depending on  $\Omega$  and  $\gamma$ .



## 1.1.2 Some Useful Inequalities

**PF** **Lemma 1.1.14.** (*Poincare-Fiedrichs Inequality*) There exists a constant  $c_0 > 0$  such that

$$\|u\|_{L^2(\Omega)} \leq c_0 \|\nabla u\|_{L^2(\Omega)} \quad \forall u \in H_0^1(\Omega) \text{ and } \forall u \in H_{\Gamma_0}^1(\Omega).$$

**Lemma 1.1.15. Cauchy-Schwartz Inequality**  $\forall f(x), g(x) \in L^2(\Omega)$  satisfies

$$\int_{\Omega} |f(x)g(x)| dx \leq \left( \int_{\Omega} |f(x)|^2 dx \right)^{\frac{1}{2}} \left( \int_{\Omega} |g(x)|^2 dx \right)^{\frac{1}{2}}$$

**Lemma 1.1.16. Hölder Inequality** Let  $1 \leq p, q \leq \infty$ . Then for all measurable real valued functions  $f$  and  $g$  satisfies the inequality

$$\int_{\Omega} |f(x)g(x)| dx \leq \left( \int_{\Omega} |f(x)|^p dx \right)^{\frac{1}{p}} \left( \int_{\Omega} |g(x)|^q dx \right)^{\frac{1}{q}},$$

where  $\frac{1}{p} + \frac{1}{q} = 1$ .

**nak1L** **Lemma 1.1.17.** For any  $a, b \in \mathbb{R}$

$$\| |a|^\gamma a - |b|^\gamma b \| \leq k_0(1 + |a|^\gamma + |b|^\gamma) |a - b|, \quad k_0, \gamma \geq 0 \quad (1.3) \quad \text{nak1}$$

*Proof.* Let us consider the function

$$f(x) = |x|^\gamma x.$$

It is clear that

$$f'(x) = (\gamma + 1)|x|^\gamma,$$

and using the mean value theorem we have

$$\begin{aligned} |f(a) - f(b)| &= |(\gamma + 1)|\theta a - (1 - \theta)b|^\gamma |(a - b)| \\ &\leq |(\gamma + 1)|(|a|^\gamma + |b|^\gamma) |(a - b)| \\ &\leq |(\gamma + 1)|(1 + |a|^\gamma + |b|^\gamma) |(a - b)|, \end{aligned}$$

where  $\theta \in (0, 1)$ .

Hence,

$$\| |a|^\gamma a - |b|^\gamma b \| \leq k_0(1 + |a|^\gamma + |b|^\gamma) |a - b|.$$

□

nak2L **Lemma 1.1.18.** For any  $x, y \in \mathbb{R}$

$$(|x|^\gamma x - |y|^\gamma y)(x - y) \geq k_1 |x - y|^{\gamma+2}, \quad k_1, \gamma \geq 0 \quad (1.4) \quad \text{nak2}$$

*Proof.*

$$\begin{aligned} J(\gamma + 2) &:= (|x|^\gamma x - |y|^\gamma y)(x - y) \\ &= \int_0^1 \frac{d}{ds} |sx + (1-s)y|^\gamma (sx + (1-s)y) ds (x - y) \\ &= \int_0^1 |sx + (1-s)y|^\gamma |x - y|^2 ds \\ &\quad + \gamma \int_0^1 |sx + (1-s)y|^{\gamma-2} ((sx + (1-s)y)(x - y))^2 ds, \end{aligned}$$

where  $0 \leq s \leq 1$ .

Since  $\gamma \geq 0$  we get;

$$J(\gamma + 2) \geq \int_0^1 |sx + (1-s)y|^\gamma |x - y|^2 ds.$$

If  $|x| \geq |y - x|$ , then

$$|sx + (1-s)y| = |x - (1-s)(x - y)| \geq |x| - (1-s)|x - y| \geq s|x - y|.$$

So we have,

$$J(\gamma + 2) \geq |x - y|^2 \int_0^1 s^\gamma |x - y|^\gamma ds = d_0 |x - y|^{\gamma+2}, \quad \text{where } d_0 = \frac{1}{\gamma + 1}$$

We showed that the inequality (1.4) is true when  $|x| \geq |y - x|$ .

If  $|x| \leq |y - x|$ , then

$$|sx + (1-s)y| = |x - (1-s)(x - y)| \leq |x| + (1-s)|x - y| \leq (2-s)|x - y|.$$

Thus we have,

$$|sx + (1-s)y|^\gamma \geq \frac{[|sx + (1-s)y|^2]^{\frac{\gamma+2}{2}}}{(2-s)|x - y|^2}.$$

Integrating the last inequality we obtain

$$\begin{aligned}
J(\gamma + 2) &\geq |x - y|^2 \int_0^1 \frac{[|sx + (1-s)y|^2]^{\frac{\gamma+2}{2}}}{(2-s)|x-y|^2} ds \geq |x|^{\gamma+2} + \frac{1}{2}|y|^{\gamma+2} \\
&\geq \frac{1}{2}(|x|^{\gamma+2} + |y|^{\gamma+2}) \geq d_0|x-y|^{\gamma+2} \geq d_0|x-y|^{\gamma+2}.
\end{aligned}$$

So we also showed that the inequality ([\(1.4\)](#)) is true when  $|x| \leq |y-x|$ .

Hence we are done.  $\square$

**Lemma 1.1.19. Young Inequality** Assume that  $x, y > 0$  and  $p, q > 1$  where  $\frac{1}{p} + \frac{1}{q} = 1$ . Then,

$$xy \leq \frac{|x|^p}{p} + \frac{|y|^q}{q}.$$

[int5L](#)

**Lemma 1.1.20. (Young Inequality with epsilon):**

For any  $x, y > 0$  and  $p, q > 1$

$$xy \leq \epsilon|x|^p + c(\epsilon)|y|^q,$$

where  $\frac{1}{p} + \frac{1}{q} = 1$  and  $c(\epsilon) = (\epsilon p)^{-\frac{q}{p}} q^{-1}$

[GenGronwall](#)

**Lemma 1.1.21. Bihari Inequality** ([\[17\]](#))

Let  $K$  be a nonnegative constant,  $\Phi$  and  $\Psi$  be non-negative continuous functions defined on  $[0, \infty)$ , and let  $g$  be a continuous non-decreasing function defined on  $[0, \infty)$  and  $g(u) > 0$  on  $(0, \infty)$ . If  $\Phi$  satisfies the following integral inequality,

$$\Phi(t) \leq K + \int_0^t \Psi(\tau)g(\Phi(\tau))d\tau$$

then

$$\Phi(t) \leq G^{-1} \left( \int_0^t \Psi(\tau)d\tau \right),$$

where  $G(\eta) = \int_K^\eta \frac{ds}{g(s)}$ ,  $\eta > K > 0$

*Proof.* Let us consider the function  $\Pi(t) : [0, \infty) \rightarrow \mathbb{R}^+$ .

$$\Pi(t) = K + \int_0^t \Psi(\tau)g(\Phi(\tau))d\tau$$

where  $K > 0$ ,  $\Phi$  and  $\Psi$  are non-negative continuous functions defined on  $[0, \infty)$ .

Then,  $\Pi$  is class of  $C^1[0, \infty)$  and

$$\Pi_t(t) = \Psi(t)g(\Phi(t)) \leq \Psi(t)g(\Pi(t)), \forall t \in [0, \infty).$$

We can apply the fundamental theorem of calculus to get

$$\frac{\Pi_t(t)}{g(\Pi(t))} = \Psi(t).$$

Integrating both sides of the inequality over the interval  $[0, t]$ ,

$$G(\Pi(t)) \leq \int_0^t \Psi(\tau) d\tau$$

Hence,

$$\Phi(t) \leq \Pi(t) \leq G^{-1}\left(\int_0^t \Psi(\tau) d\tau\right),$$

□

ClasGron **Lemma 1.1.22. (Gronwall Inequality):** <sup>Nis2</sup> ~~[7]~~ If  $g(s) = Cs$  then

$$\Phi(t) \leq K + C \int_0^t \Psi(\tau)\Phi(\tau) d\tau$$

implies

$$\Phi(t) \leq K \exp\left(C \int_0^t \Psi(\tau) d\tau\right).$$

*Proof.* Taking  $g(s) = Cs$  in Lemma <sup>GenGronwall</sup> 1.1.21, then  $G^{-1}(\eta) = K \exp(\eta)$  and we obtain

$$\Phi(t) \leq K \exp\left(C \int_0^t \Psi(\tau) d\tau\right).$$

□

Int **Lemma 1.1.23.** For each  $v \in H^2(\Omega) \cap H_0^1(\Omega)$  the following inequality holds true

$$\|\nabla u\|^2 \leq \|u\| \|\Delta u\| \tag{1.5} \span style="border: 1px solid black; padding: 2px;">inter1$$

*Proof.* Since  $u = 0$  on the boundary  $\partial\Omega$  we have

$$0 = \int_{\Omega} \nabla \cdot (u \nabla u) dx = \|\nabla u\|^2 + (u, \Delta u).$$

Hence due to the Cauchy-Schwarz inequality we get

$$\|\nabla u\|^2 = -(u, \Delta u) \leq \|u\| \|\Delta u\|.$$

□

### 1.1.3 Auxiliary material

**kom1T** **Theorem 1.1.24.** <sup>kom1</sup>~~(1.3)~~ Let  $E := \mathbb{R}^+ \rightarrow \mathbb{R}^+$  be a nonincreasing function, and assume that there exists a constant  $T > 0$  such that

$$\int_t^{\infty} E(s) ds \leq TE(t) \quad \forall t \in \mathbb{R}^+ \quad (1.6) \quad \text{kom2}$$

then

$$E(t) \leq E(0)e^{\frac{1-t}{T}} \quad \forall t \geq T. \quad (1.7) \quad \text{kom3}$$

*Proof.* Let us consider the following function

$$f(x) := e^{\frac{x}{T}} \int_x^{\infty} E(s) ds, \quad x \in \mathbb{R}^+.$$

The function  $f$  is locally absolutely continuous and nonincreasing on  $\mathbb{R}^+$ . In fact, by using the fundamental theorem of calculus and the condition <sup>kom2</sup>~~(1.6)~~ we get

$$f'(x) = \frac{1}{T} e^{\frac{x}{T}} \int_x^{\infty} E(s) ds + e^{\frac{x}{T}} (-E(x)) = \frac{1}{T} e^{\frac{x}{T}} \left( \int_x^{\infty} E(s) ds - TE(x) \right) \leq 0, \quad \forall x \in \mathbb{R}^+$$

So  $f$  is nonincreasing. Therefore  $f(x) \leq f(0)$ ,  $\forall x \in \mathbb{R}^+$ , and due to <sup>kom2</sup>~~(1.6)~~ we get,

$$\int_x^{\infty} E(s) ds = f(x) \leq f(0) = \int_0^{\infty} E(s) ds \leq TE(0), \quad \forall x \in \mathbb{R}^+.$$

So we have

$$\int_x^{\infty} E(s) ds \leq TE(0)e^{-\frac{x}{T}}, \quad \forall x \in \mathbb{R}^+. \quad (1.8) \quad \text{kom4}$$

Since  $E$  is nonnegative and nonincreasing, we have

$$\int_x^\infty E(s)ds \geq \int_x^{x+T} E(s)ds \geq (x+T-x)E(x+T) = TE(x+T). \quad (1.9) \quad \boxed{\text{kom5}}$$

By  $\boxed{\text{kom4}}$  and  $\boxed{\text{kom5}}$

$$E(x+T) \leq E(0)e^{-\frac{x}{T}} \quad \forall x \in \mathbb{R}^+.$$

Thus we have

$$E(t) \leq E(0)e^{\frac{1-t}{T}},$$

where  $t := x + T$ . □

$\boxed{\text{nak4L}}$  **Lemma 1.1.25.**  $\boxed{\text{nak}}$   $\boxed{\text{oj}}$  Let  $\Phi(t)$  be a positive function on  $\mathbb{R}^+$  satisfying :

$$k\Phi(t)^{\alpha+1} \leq \Phi(t) - \Phi(t+1) \quad \forall t \geq 0, \quad (1.10) \quad \boxed{\text{n}}$$

for some constants  $k$  and  $\alpha \geq 0$ . Then we have

$$\Phi(t) \leq (\alpha k(t-1) + M^{-\alpha})^{-\frac{1}{\alpha}}, \quad \forall t \geq 0$$

where  $M = \max_{t \in [0,1]} \Phi(t)$

*Proof.* Let us denote  $\Phi(t)^{-\alpha} = y(t)$ .

$$\begin{aligned} y(t+1) - y(t) &= \int_0^1 \frac{d}{d\theta} (\theta\Phi(t+1) + (1-\theta)\Phi(t))^{-\alpha} d\theta \\ &= \int_0^1 -\alpha(\theta\Phi(t+1) + (1-\theta)\Phi(t))^{-\alpha-1} (\Phi(t+1) - \Phi(t)) d\theta \\ &\geq \alpha k \Phi(t)^{\alpha+1} \int_0^1 \Phi(t)^{-\alpha-1} d\theta = \alpha k. \end{aligned}$$

For  $\forall t \geq 1$ , choose the integer  $n$  as  $n \leq t \leq n+1$  and by  $\boxed{\text{I.10}}$  we get,

$$y(t) - y(t-n) \geq n\alpha k.$$

Since

$$nk\Phi(t)^{\alpha+1} \leq \Phi(t) - \Phi(t-n),$$

we have

$$y(t) \geq y(t-n) + n\alpha k \geq y(t-n) + (t-1)\alpha k$$

$$\Phi(t)^{-\alpha} \geq (t-1)\alpha k + \Phi(t-n)^{-\alpha}.$$

Then

$$\Phi(t) \leq ((t-1)\alpha k + \Phi(t-n)^{-\alpha})^{\frac{-1}{\alpha}} \leq ((t-1)\alpha k + \max_{t \in [0,1]} \Phi(t)^{-\alpha})^{\frac{-1}{\alpha}} = ((t-1)\alpha k + M^{-\alpha})^{\frac{-1}{\alpha}}.$$

□

**nak5L** **Lemma 1.1.26.** <sup>nak</sup><sub>[5]</sub> Let  $\Phi(t)$  be a positive function that satisfies the inequality:

$$k\Phi(t) \leq \Phi(t) - \Phi(t+1) \quad \text{for } \forall t \geq 0. \quad (1.11) \quad \text{nak4}$$

with some constant  $k < 1$ .

Then the following estimate holds true:

$$\Phi(t) \leq Me^{-k't}, \quad t \geq 1. \quad (1.12) \quad \text{nak5}$$

Here  $k' = -\log(1-k) > 0$

*Proof.* By <sup>nak4</sup><sub>(I.11)</sub>

$$\Phi(t+1) \leq (1-k)\Phi(t).$$

Therefore, if  $t \geq 1$ , for  $n$  such that  $n \leq t \leq n+1$ , we have

$$\Phi(t) \leq (1-k)\Phi(t-1).$$

Since  $0 < k < 1$ ,  $\frac{1}{1-k} \geq 1-k$  and  $\Phi(t) \leq \frac{1}{1-k}\Phi(t-1)$ .

Thus,

$$\Phi(t) \leq (1-k)^2\Phi(t-2) \cdots \leq (1-k)^n\Phi(t-n) \leq \frac{1}{(1-k)^n}\Phi(t-n)$$

$$\Phi(t) \leq \left(\frac{1}{1-k}\right)^n\Phi(t-n)$$

$$\Phi(t) \leq \left(\frac{1}{1-k}\right)^n\Phi(t-n) \leq M(1-k)^{-t} = Me^{-t\log(1-k)} = Me^{-k't}.$$

□

nis4L **Lemma 1.1.27.** *Let  $\phi(t)$  be a bounded nonnegative function on  $\mathbb{R}^+$ , satisfying*

$$\max_{t \leq \tau \leq t+1} \phi(\tau) \leq K_0(\phi(t) - \phi(t+1)) + g(t) \quad (1.13) \quad \text{int1}$$

where  $K_0$  is a positive constant and  $g(t)$  is a nonnegative function such that

$$g(t) \leq C_0 \exp(-\theta t) \text{ with } \theta > 0, C_0 > 0 \quad (1.14) \quad \text{int2}$$

then

$$\phi(t) \leq C' \exp(-\theta' t) \quad (1.15) \quad \text{int3}$$

where  $\theta' = \min(-\theta, \log \frac{K_0}{K_0+1})$  and  $C'$  is a positive constant depending on  $K_0, C_0, \theta$ .

*Proof.* By int1 (1.13),

$$\phi(t+1) \leq K_0(\phi(t) - \phi(t+1)) + g(t)$$

then

$$\phi(t+1) \leq \frac{K_0}{K_0+1} \phi(t) + \frac{K_0}{K_0+1} g(t)$$

If  $t \geq 1$ , there exists  $n$  such that  $n \leq t \leq n+1$ . Then we have

$$\phi(t) \leq \left( \frac{K_0}{K_0+1} \right)^n \phi(t-n) + \sum_{i=1}^n \left( \frac{K_0}{K_0+1} \right)^i g(t)$$

Since  $\frac{K_0}{K_0+1} < 1$ , then

$$\begin{aligned} \phi(t) &\leq \left( \frac{K_0}{K_0+1} \right)^n \phi(t-n) + \sum_{i=1}^n \left( \frac{K_0}{K_0+1} \right)^i g(t) \\ &\leq M \left( \frac{K_0}{K_0+1} \right)^t + \sum_{i=1}^n \left( \frac{K_0}{K_0+1} \right)^i g(t) \\ &\leq M \exp \left( \log \left( \frac{K_0}{K_0+1} \right) \right) + C_1 \exp(-\theta t) \leq C' \exp(-\theta' t) \end{aligned}$$

where  $\theta' = \min(-\theta, \log \frac{K_0}{K_0+1})$  and  $M = \max_{t \in [0,1]} \Phi(t)$ .

Hence we are done. □



## Chapter 2

# Estimates of Solutions of Linear Wave Equation

The aim of this chapter is to show the exponential energy decay of the solutions of the wave equation under suitable linear boundary conditions. We use the Theorem <sup>kom11</sup> I.I.24 to show that energy norm of the solution of the problem satisfies this theorem. Hence, we conclude that the problem has energy decay. In other words, it is asymptotically stable.

### 2.1 Linear wave equation

In this section we consider the following problem:

$$u_{tt} - \Delta u + qu = 0, \quad x \in \Omega, t \in \mathbb{R}^+ \quad (2.1) \quad \boxed{\text{lw1}}$$

$$u(x, t) = 0, \quad x \in \Gamma_0, t \in \mathbb{R}^+ \quad (2.2) \quad \boxed{\text{lw2}}$$

$$\partial_\nu u + au + lu_t = 0, \quad x \in \Gamma_1, t \in \mathbb{R}^+ \quad (2.3) \quad \boxed{\text{lw3}}$$

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), x \in \Omega, \quad (2.4) \quad \boxed{\text{lw42}}$$

where  $\Omega \subset \mathbb{R}^n$  ( $n \geq 3$ ) is a bounded domain with a  $\mathbf{C}^2$  boundary  $\Gamma$  and  $\{\Gamma_0, \Gamma_1\}$  is a partition of boundary  $\Gamma, q : \Omega \rightarrow \mathbb{R}, a, l : \Gamma_1 \rightarrow \mathbb{R}^+$  are

continuous and non negative functions. We assume that the initial data  $u_0$  and  $u_1$  are sufficiently smooth functions so that the problem <sup>lw1</sup> (2.1)-<sup>lw42</sup> (2.4) has a classical solution.

The following lemma shows that the energy is a non-increasing function of time.

kom1L **Lemma 2.1.1.** <sup>[Kom1]</sup> <sup>[3]</sup> Suppose that the problem <sup>[w1]</sup><sub>(2.1)</sub>-<sup>[w42]</sup><sub>(2.4)</sub> has a classical solution. Then the solution of the problem <sup>[w1]</sup><sub>(2.1)</sub>-<sup>[w42]</sup><sub>(2.4)</sub> satisfies the energy equality:

$$E(S) - E(T) = \int_S^T \int_{\Gamma_1} l(u_t(t))^2 d\Gamma, \quad \text{such that } 0 \leq S < T < \infty. \quad (2.5) \quad \text{kom1}$$

*Proof.* We multiply the equation <sup>[w1]</sup><sub>(2.1)</sub> by  $u_t$  and integrate over the region  $\Omega \times (S, T)$ :

$$0 = \int_S^T \int_{\Omega} \frac{1}{2} ((u_t)^2)_t dx dt - \int_S^T \int_{\Omega} u_t \Delta u dx dt + \int_S^T \int_{\Omega} \frac{1}{2} q(u^2)_t dx dt.$$

By using the equality  $\nabla u_t \nabla u = \frac{1}{2} (|\nabla u|^2)_t$  and Green's identity we get

$$0 = \frac{1}{2} \left( \int_{\Omega} (u_t)^2 + qu^2 + |\nabla u|^2 dx \right) \Big|_S^T - \int_S^T \int_{\Gamma} u_t d_{\nu} u dx dt.$$

Since  $\partial_{\nu} u + au + lu_t = 0$ , we have

$$0 = \frac{1}{2} \left( \int_{\Omega} (u_t)^2 + qu^2 + |\nabla u|^2 dx + \int_{\Gamma_1} au^2 d\Gamma \right) \Big|_S^T + \int_S^T \int_{\Gamma_1} l(u_t)^2 dx dt$$

or

$$\int_S^T \int_{\Gamma_1} l(u_t)^2 dx dt = E(S) - E(T).$$

Hence  $E(S) \geq E(T)$ . □

It follows from Lemma <sup>[kom1L]</sup><sub>2.1.1</sub> that  $lu_t \in L^2(\mathbb{R}^+, L^2(\Gamma_1))$ .

In this section our purpose is to show that the energy function tends to zero with an exponential rate as  $t \rightarrow \infty$  for particular choice of functions  $l$  and  $a$ .

Assume that there is a point  $x_0 \in \mathbb{R}^n$  such that  $m \cdot \nu \leq 0$  on  $\Gamma_0$  and  $m \cdot \nu \geq 0$  on  $\Gamma_1$ , where  $m(x) = x - x_0$  and  $\lambda_1$  is the biggest constant for which the following inequality holds:

$$\int_{\Omega} |\nabla u|^2 + qu^2 dx + \int_{\Gamma_1} au^2 d\Gamma \geq \lambda_1 \int_{\Omega} u^2 dx. \quad (2.6) \quad \text{kom7}$$

By Poincare-Friedrich's inequality and nonnegativity of the functions  $a$  and  $q$  such a number  $\lambda_1$  exists.

In what follows we use the notations:

$$k := \frac{1}{R}, \quad b := \frac{(n-1)}{2R^2}, \quad Mu := 2m \cdot \nabla u + (n-1)u, \quad \Gamma_m = (m \cdot \nu) d\Gamma$$

$$l := \frac{(m \cdot \nu)}{R}, \quad a := \frac{(n-1)(m \cdot \nu)}{2R^2}$$

kom2L **Lemma 2.1.2.** <sup>[Kom1]</sup><sub>([3])</sub> For any given  $(u_0, u_1)$  and  $0 \leq S \leq T < \infty$  arbitrarily, the solution of the problem satisfies the following inequality:

$$2 \int_S^T E(t) dt \leq \left[ \int_{\Omega} u_t M u dx \right]_S^T - \int_S^T \int_{\Omega} (n-2) q u^2 + 2 q u m \cdot \nabla u dx dt$$

$$+ \int_S^T \int_{\Gamma_0} |\partial_{\nu} u|^2 d\Gamma_m dt$$

$$+ \int_S^T \int_{\Gamma_1} u_t^2 - |\nabla u|^2 + b u^2 - (k u_t + b u) M u d\Gamma_m dt. \quad (2.7) \quad \text{kom10}$$

*Proof.* We multiply <sup>[Lw1]</sup><sub>(2.1)</sub> by  $2m \cdot \nabla u$  and integrate over the region  $\Omega \times (S, T)$  use the Green's identity and the identities

$$2u_{tt}m \cdot \nabla u = (2u_t m \cdot \nabla u)_t - 2u_t m \cdot \nabla u_t, \quad 2u_t \nabla u_t = \nabla u_t^2$$

we get,

$$\int_S^T \int_{\Omega} -q u 2m \cdot \nabla u dx dt = \left( \int_{\Omega} 2u_t m \cdot \nabla u \right) \Big|_S^T$$

$$+ \int_S^T \int_{\Omega} 2\nabla u \cdot \nabla (m \cdot \nabla u) - m \cdot \nabla u_t^2 dx dt$$

$$- \int_S^T \int_{\Gamma} 2\partial_{\nu} u m \cdot \nabla u d\Gamma dt.$$

Now, applying divergence theorem and using the identities

$$-m \cdot \nabla u_t^2 = -\nabla \cdot (m \cdot u_t^2) + \nabla \cdot m (u_t^2), \quad 2\nabla u \cdot \nabla (m \cdot \nabla u)$$

$$= 2 \sum_{i,j=1}^n \partial_i u \partial_i (m_j \partial_j u), \quad \partial_j ((\partial_i u)^2) = 2\partial_i \partial_j u \partial_i u$$

we get,

$$\begin{aligned}
\int_S^T \int_{\Omega} -qu2m \cdot \nabla u dx dt &= \left( \int_{\Omega} 2u_t m \cdot \nabla u dx \right) \Big|_S^T \\
&\quad - \int_S^T \int_{\Gamma} 2\partial_{\nu} u m \cdot \nabla u + (m \cdot \nu) u_t^2 d\Gamma dt \\
&\quad + \int_S^T \int_{\Omega} (\nabla \cdot m) u_t^2 + 2 \sum_{i,j=1}^n \partial_i m_j \partial_i u \partial_j u + \sum_{i,j=1}^n m_j \partial_j ((\partial_i u)^2) dx dt. \quad (2.8) \quad \boxed{\text{kom43}}
\end{aligned}$$

Again by divergence theorem and the identities

$$\begin{aligned}
m \nabla (|\nabla u|^2) &= \nabla (m |\nabla u|^2) - (\nabla \cdot m) |\nabla u|^2, \\
\sum_{i,j=1}^n \partial_i m_j \partial_i u \partial_j u &= 2|\nabla u|^2, \quad \nabla \cdot m = n
\end{aligned}$$

the equality  $\stackrel{\boxed{\text{kom43}}}{(2.8)}$  can be written as,

$$\begin{aligned}
\int_S^T \int_{\Gamma} 2\partial_{\nu} u m \cdot \nabla u + (m \cdot \nu) (u_t^2 - |\nabla u|^2) d\Gamma dt &= \\
\left( \int_{\Omega} 2u_t m \cdot \nabla u dx \right) \Big|_S^T + \int_S^T \int_{\Omega} n u_t^2 - (2-n) |\nabla u|^2 + 2qu m \nabla u dx dt. \quad (2.9) \quad \boxed{\text{kom11}}
\end{aligned}$$

Let us multiply the equation  $\stackrel{\boxed{\text{1w1}}}{(2.9)}$  by  $(n-1)u$ , integrate over  $\Omega \times (S; T)$  and use the Green's identity:

$$\begin{aligned}
0 &= (n-1) \int_{\Omega} u_t u \Big|_S^T + (n-1) \int_S^T \int_{\Omega} -u_t^2 + |\nabla u|^2 + qu^2 dx dt \\
&\quad - (n-1) \int_S^T \int_{\Gamma} u \partial_{\nu} u d\Gamma dt. \quad (2.10) \quad \boxed{\text{kom12}}
\end{aligned}$$

Adding  $\stackrel{\boxed{\text{kom11}}}{(2.9)}$  and  $\stackrel{\boxed{\text{kom12}}}{(2.10)}$ , we get

$$\begin{aligned}
&\int_S^T \int_{\Gamma} \partial_{\nu} u M u + (m \cdot \nu) (u_t^2 - |\nabla u|^2) d\Gamma dt \\
&= \left[ \int_{\Omega} u_t M u \right] \Big|_S^T + \int_S^T \int_{\Omega} u_t^2 + |\nabla u|^2 + (n-1)qu^2 + 2qu m \cdot \nabla u dx dt. \quad (2.11) \quad \boxed{\text{kom45}}
\end{aligned}$$

Let us rewrite  $\frac{\text{kom45}}{(2.11)}$  in the form,

$$\begin{aligned}
& \int_S^T \int_{\Gamma} \partial_{\nu} u M u + (m \cdot \nu)(u_t^2 - |\nabla u|^2) d\Gamma dt \\
&= \left[ \int_{\Omega} u_t M u \right]_S^T + \int_S^T \int_{\Omega} u_t^2 + |\nabla u|^2 + q u^2 + (n-2) q u^2 + 2 q u m \cdot \nabla u dx dt \\
&\quad + \int_S^T \int_{\Gamma_1} \frac{(n-1)(m \cdot \nu)}{2R^2} u^2 d\Gamma dt - \int_S^T \int_{\Gamma_1} \frac{(n-1)(m \cdot \nu)}{2R^2} u^2 d\Gamma dt. \quad (2.12) \quad \boxed{\text{kom46}}
\end{aligned}$$

i.e.

$$\begin{aligned}
& \int_S^T \int_{\Gamma} \partial_{\nu} u M u + (m \cdot \nu)(u_t^2 - |\nabla u|^2) d\Gamma dt \\
&= \left[ \int_{\Omega} u_t M u \right]_S^T + \int_S^T E(t) dt \\
&\quad + \int_S^T \int_{\Omega} (n-2) q u^2 + 2 q u m \cdot \nabla u dx dt - \int_S^T \int_{\Gamma_1} b u^2 d\Gamma_m dt. \quad (2.13) \quad \boxed{\text{kom55}}
\end{aligned}$$

Since  $u = 0$  on  $\Gamma_0$  we have  $\nabla u = \partial_{\nu} u \cdot \nu$  on  $\Gamma_0$ . Thus,

$$\int_{\Gamma_0} \partial_{\nu} u M u d\Gamma = \int_{\Gamma_0} \partial_{\nu} u (2m \cdot \nabla u) d\Gamma = 2 \int_{\Gamma_0} (\partial_{\nu} u)^2 (m \cdot \nu) d\Gamma, \quad (2.14) \quad \boxed{\text{kom13}}$$

$$\int_{\Gamma_0} -|\nabla u|^2 (m \cdot \nu) d\Gamma = - \int_{\Gamma_0} (\partial_{\nu} u)^2 |\nu|^2 (m \cdot \nu) d\Gamma = - \int_{\Gamma_0} (\partial_{\nu} u)^2 (m \cdot \nu) d\Gamma. \quad (2.15) \quad \boxed{\text{kom14}}$$

Therefore by  $\frac{\text{kom13}}{(2.14)}$  and  $\frac{\text{kom14}}{(2.15)}$  the left hand side of  $\frac{\text{kom55}}{(2.13)}$  can be written

as :

$$\begin{aligned}
& \int_S^T \int_{\Gamma} \partial_{\nu} u M u + (m \cdot \nu)(u_t^2 - |\nabla u|^2) d\Gamma dt \\
&= \int_S^T \int_{\Gamma} \partial_{\nu} u M u + \int_S^T \int_{\Gamma_1} (u_t^2 - |\nabla u|^2) d\Gamma_m dt + \int_S^T \int_{\Gamma_0} -|\nabla u|^2 d\Gamma_m dt \\
&= \int_S^T \int_{\Gamma_1} (-l u_t - a u) M u d\Gamma dt + 2 \int_S^T \int_{\Gamma_0} (\partial_{\nu} u)^2 d\Gamma_m dt \\
&\quad + \int_S^T \int_{\Gamma_1} (u_t^2 - |\nabla u|^2) d\Gamma_m dt - \int_S^T \int_{\Gamma_0} (\partial_{\nu} u)^2 d\Gamma_m dt \\
&= \int_S^T \int_{\Gamma_0} (\partial_{\nu} u)^2 d\Gamma_m dt + \int_S^T \int_{\Gamma_1} u_t^2 - |\nabla u|^2 - (k u_t + b u) M u d\Gamma_m dt.
\end{aligned}$$

So we get the inequality:

$$\begin{aligned}
2 \int_S^T E(t) dt &\leq \left[ \int_{\Omega} u_t M u dx \right]_S^T - \int_S^T \int_{\Omega} (n-2) q u^2 + 2 q u m \cdot \nabla u dx dt \\
&\quad + \int_S^T \int_{\Gamma_0} |\partial_{\nu} u|^2 d\Gamma_m dt + \int_S^T \int_{\Gamma_1} u_t^2 - |\nabla u|^2 + b u^2 - (k u_t + b u) M u d\Gamma_m dt.
\end{aligned}$$

Proof follows. □

kom3L **Lemma 2.1.3.** <sup>[Kom2]</sup><sub>([4])</sub> Let  $u$  be solution of the problem <sup>[1w1]</sup><sub>(2.1)</sub>-<sup>[1w42]</sup><sub>(2.4)</sub>. Then the following estimate holds:

$$\left| \int_{\Omega} u_t M u dx \right| \leq 2 R E(t), \quad \forall t \in \mathbb{R}^+. \quad (2.16) \quad \text{kom15}$$

*Proof.* By Young's inequality with epsilon

$$\left| \int_{\Omega} u_t M u dx \right| \leq \int_{\Omega} R |u_t|^2 + \frac{|M u|^2}{4R} dx.$$

$$\begin{aligned}\int_{\Omega} |Mu|^2 dx &= \int_{\Omega} |2m \cdot \nabla u|^2 + (n-1)^2 u^2 + 4(n-1)um \cdot \nabla u dx \\ &= \int_{\Omega} |2m \cdot \nabla u|^2 + (n-1)^2 u^2 + (2n-2)m \cdot 2u \nabla u dx.\end{aligned}$$

We use divergence theorem and the identities,

$$2u \nabla u = \nabla u^2,$$

$$(2n-2)m \cdot \nabla u^2 = -(2n-2)nu^2 + (2n-2)\nabla \cdot (mu^2)$$

to get

$$\begin{aligned}\int_{\Omega} |Mu|^2 dx &= \int_{\Omega} |2m \cdot \nabla u|^2 + ((n-1)^2 - (2n-2)n)u^2 dx \\ &\quad + (2n-2) \int_{\Gamma} (m \cdot \nu)u^2 d\Gamma = \int_{\Omega} |2m \cdot \nabla u|^2 + (1-n^2)u^2 dx \\ &\quad + (2n-2) \int_{\Gamma} (m \cdot \nu)u^2 d\Gamma. \quad (2.17) \quad \boxed{\text{kom56}}\end{aligned}$$

Since  $(1-n^2)u^2$  is negative, then we can majorize  $\overset{\boxed{\text{kom56}}}{(2.17)}$  as:

$$\begin{aligned}\int_{\Omega} |Mu|^2 dx &\leq 4R^2 \int_{\Omega} |\nabla u|^2 dx + 4R^2 \frac{n-1}{2R^2} \int_{\Gamma} (m \cdot \nu)u^2 d\Gamma \\ &\leq 4R^2 \int_{\Omega} |\nabla u|^2 dx + 4R^2 \int_{\Gamma} au^2 d\Gamma.\end{aligned}$$

As a result:

$$\int_{\Omega} |u_t Mu| dx \leq R \int_{\Omega} |u_t|^2 dx + \frac{4R^2}{4R} \left( \int_{\Omega} |\nabla u|^2 + qu^2 dx + \int_{\Gamma} au^2 d\Gamma \right) = 2RE.$$

□

$\boxed{\text{kom2T}}$  **Theorem 2.1.4.**  $\overset{\boxed{\text{Kom1}}}{(1.3)}$  Suppose that

- $q \in C(\bar{\Omega})$ ,
- $Q_1 = 2R \sup_{\Omega} q / \sqrt{\lambda_1} < 1$  if  $n \geq 2$ ,

Then for every given  $(u_0, u_1)$  the solution of the problem satisfies

$$E(t) \leq E(0) \exp\left(\frac{1 - (1 - Q_1)t}{2R}\right), \forall t \in \mathbb{R}^+ \quad (2.18) \quad \boxed{\text{kom9}}$$

*Proof.* To use the Theorem [2.1.4](#) <sup>kom2T</sup> we estimate the terms on the right hand side of the inequality [\(2.7\)](#) <sup>kom10</sup>.

By Lemma [2.1.3](#) <sup>kom3L</sup>, the first term of the right hand side of the inequality [\(2.7\)](#) <sup>kom10</sup> can be estimated as:

$$\left[ \int_{\Omega} u_t M u dx \right]_S^T \leq 2RE(t)|_S^T \leq 2RE(S) + 2RE(T).$$

Since  $m \cdot \nu \leq 0$  on  $\Gamma_0$  <sup>kom10</sup>, we can ignore the third term of the right hand side of the inequality [\(2.7\)](#) <sup>kom10</sup>. i.e.

$$\int_S^T \int_{\Gamma_0} |\partial_\nu u|^2 d\Gamma_m dt = \int_S^T \int_{\Gamma_0} (m \cdot \nu) |\partial_\nu u|^2 d\Gamma dt \leq 0$$

For the second term we will show that:

$$- \int_S^T \int_{\Omega} (n-2)qu^2 + 2qum \cdot \nabla u dx dt \leq 2Q_1 \int_S^T E(t) dt.$$

Since  $n \geq 3$  and  $q$  is a positive function then

$$- \int_S^T \int_{\Omega} (n-2)qu^2 dx dt \leq 0.$$

By using Young's inequality we get,

$$|u \nabla u| \leq \frac{\lambda_1^{\frac{1}{2}}}{2} u^2 + \frac{1}{2\lambda_1^{\frac{1}{2}}} |\nabla u|^2. \quad (2.19) \quad \boxed{\text{kom16}}$$



Therefore by [\(2.6\)](#) and [\(2.19\)](#), we have,

$$\begin{aligned}
-\int_S^T \int_{\Omega} 2qum \cdot \nabla u dx dt &\leq 2R \sup_{\Omega} q \int_S^T \int_{\Omega} |u \nabla u| dx dt \\
&\leq 2R \sup_{\Omega} q \int_S^T \frac{\lambda_1^{\frac{1}{2}}}{2} \left( \int_{\Omega} u^2 dx \right) + \frac{1}{2\lambda_1^{\frac{1}{2}}} \left( \int_{\Omega} |\nabla u|^2 dx \right) dt \\
&\leq 2R \sup_{\Omega} q \int_S^T \frac{1}{2\lambda_1^{\frac{1}{2}}} \left( \int_{\Omega} |\nabla u|^2 + qu^2 dx + \int_{\Gamma_1} au^2 d\Gamma \right) + \frac{1}{2\lambda_1^{\frac{1}{2}}} \left( \int_{\Omega} |\nabla u|^2 dx \right) dt \\
&\leq 2R \sup_{\Omega} q \lambda_1^{-\frac{1}{2}} \int_S^T \left( \int_{\Omega} u_t^2 + |\nabla u|^2 + qu^2 dx + \int_{\Gamma_1} au^2 d\Gamma \right) dt \\
&\leq 2Q_1 \int_S^T E(t) dt.
\end{aligned}$$

Thus, we can deduce that:

$$\begin{aligned}
2(1 - Q_1) \int_S^T E(t) dt &\leq 2RE(S) + 2RE(T) \\
&\quad + \int_S^T \int_{\Gamma_1} u_t^2 - |\nabla u|^2 + bu^2 - (ku_t + bu)Mu d\Gamma_m dt.
\end{aligned}$$

By Young's inequality, we obtain

$$2(ku_t + bu)m \cdot \nabla u \leq R^2(ku_t + bu)^2 + |\nabla u|^2. \quad (2.20) \quad \boxed{\text{kom57}}$$

We use the particular choice of  $k$ ,  $b$  and [\(2.20\)](#) to make the following estimation on the last term of the right hand side of the inequality [\(2.7\)](#). [\(kom10\)](#)

$$\begin{aligned}
&u_t^2 - |\nabla u|^2 + bu^2 - (ku_t + bu)Mu \\
&= u_t^2 - |\nabla u|^2 + bu^2 - 2(ku_t + bu)m \cdot \nabla u + (1 - n)u(ku_t + bu) \\
&\leq u_t^2 - |\nabla u|^2 + bu^2 + R^2(ku_t + bu)^2 + |\nabla u|^2 + (1 - n)u(ku_t + bu) \\
&= u_t^2 + b(2 - n + R^2b)u^2 + R^2k^2(u_t)^2 + (1 - n + 2R^2b)ku_t u \\
&= 2u_t^2 + ((3 - n)/2)bu^2.
\end{aligned}$$

So we deduce that,

$$2(1 - Q_1) \int_S^T E(t) dt \leq 2RE(S) + 2RE(T) + 2 \int_S^T \int_{\Gamma_1} u_t^2 d\Gamma_m dt + ((3 - n)/2) \int_S^T \int_{\Gamma_1} bu^2 d\Gamma_m dt. \quad (2.21) \quad \boxed{\text{kom50}}$$

According to the particular choice of  $l$ ,

$$E(S) - E(T) = \int_S^T \int_{\Gamma_1} l(u_t)^2 d\Gamma dt = \frac{1}{R} \int_S^T \int_{\Gamma_1} (u_t)^2 d\Gamma_m dt.$$

Then, the third term of  $(\text{kom50})$  will be:

$$2RE(S) - 2RE(T) = 2 \int_S^T \int_{\Gamma_1} (u_t)^2 d\Gamma_m dt.$$

Thus,  $(\text{kom50})$  can be majorized as:

$$2(1 - Q_1) \int_S^T E(t) dt \leq 4RE(S) + ((3 - n)/2) \int_S^T \int_{\Gamma_1} bu^2 d\Gamma_m dt. \quad (2.22) \quad \boxed{\text{kom51}}$$

Since  $n \geq 3$  last term of the  $(\text{kom51})$  is less than 0, then we get

$$2(1 - Q_1) \int_S^T E(t) dt \leq 4RE(S).$$

Let  $T$  tend to infinity and  $S \in \mathbb{R}^+$  is fixed, then we obtain,

$$\int_S^\infty E(t) dt \leq \frac{2R}{(1 - Q_1)} E(S).$$

Finally, we apply theorem  $(\text{kom1T})$  and conclude that

$$E(t) \leq E(0) \exp\left(1 - \frac{(1 - Q_1)t}{2R}\right), \forall t \in \mathbb{R}^+.$$

□

Hence, by Theorem  $(\text{kom2T})$ , we show the asymptotic stability of the problem.

## Chapter 3

# Estimates of Solutions of Wave Equation with Nonlinear Damping Term

In this chapter we consider the wave equation with the nonlinear damping term. We obtain a polynomial decay estimate for the difference of two solutions of the equation under consideration.

### 3.1 Wave Equation with Nonlinear Damping Term

The problem that we consider in this chapter is:

$$u_{tt}(x, t) - \Delta u(x, t) + |u_t(x, t)|^\gamma u_t(x, t) = f(x, t), \quad x \in \Omega, t > 0 \quad (3.1) \quad \boxed{\text{nk1}}$$

$$u(x, t) = 0, \quad x \in \partial\Omega, t > 0 \quad (3.2) \quad \boxed{\text{nk2}}$$

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad x \in \Omega, \quad (3.3) \quad \boxed{\text{nk3}}$$

where  $\Omega \subset \mathbb{R}^n$  is a bounded domain with sufficiently smooth boundary  $\partial\Omega$  and  $u_0(x)$ ,  $u_1(x)$  and  $f(x, t)$  are sufficiently smooth functions then the problem (3.1)-(3.3) has classical solutions.

We suppose that:

$$0 \leq \gamma \leq \frac{4}{n-2} \text{ if } n \geq 3, \quad 0 \leq \gamma < \infty \text{ if } n = 1, 2.$$

$$\|u(t)\|_E^2 = \|u_t(t)\|_{L^2(\Omega)}^2 + \|u(t)\|_{H_0^1(\Omega)}^2$$

$$\|f(t)\|_S = \left( \int_t^{t+1} \int_\Omega f(x, s)^{\frac{\gamma+2}{\gamma+1}} dx ds \right)^{\frac{\gamma+1}{\gamma+2}}$$

nk3T

**Theorem 3.1.1.** <sup>nak</sup> [5] Let  $u(t), v(t)$  be any two solution of the problem <sup>nk1</sup> (3.1)-  
<sup>nk3</sup> (3.3) then  $u, v$  satisfies

$$\|u(t) - v(t)\|_E \leq \left(\frac{\gamma}{2}K(t-1) + M^{-\gamma}\right)^{-\frac{1}{\gamma}} \text{ for } t \geq 1 \text{ if } \gamma > 0$$

and

$$\|u(t) - v(t)\|_E \leq Me^{-K't} \text{ for } t \geq 1 \text{ if } \gamma = 0$$

where  $K' = -\frac{1}{2}\log(1-K) > 0$  and  $M = \max_{t \in [0,1]} \|u(t) - v(t)\|_E$   
and  $K$  is positive constant.

*Proof.* We multiply the equation <sup>nk1</sup> (3.1) by  $u_t(x, t)$ , integrate over  $\Omega \times [t, t+1]$   
and get

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_t^{t+1} \int_{\Omega} u_t^2(x, s) dx ds + \frac{1}{2} \frac{d}{dt} \int_t^{t+1} \int_{\Omega} |\nabla u(x, s)|^2 dx ds + \int_t^{t+1} \int_{\Omega} |u_t(x, s)|^{\gamma+2} dx ds \\ = \int_t^{t+1} \int_{\Omega} u_t(x, s) f(x, s) dx ds. \end{aligned}$$

Taking  $k_2 \leq 1$ , we have

$$\begin{aligned} k_2 \int_t^{t+1} \int_{\Omega} |u_t(x, s)|^{\gamma+2} dx ds \leq \\ \|u(t)\|_E^2 - \|u(t+1)\|_E^2 + \int_t^{t+1} \int_{\Omega} u_t(x, s) f(x, s) dx ds. \end{aligned} \quad (3.4) \quad \text{nak60}$$

If  $p = \gamma + 2$ ,  $q = \frac{\gamma+2}{\gamma+1}$ ,  $\epsilon = \frac{k_2}{2}$ , by Young's inequality with epsilon we get

$$u_t(x, t) f(x, t) \leq \frac{k_2}{2} |u_t(x, t)|^{\gamma+2} + \frac{\gamma+2}{\gamma+1} \left(\frac{2}{k_2\gamma+2}\right) |f(x, t)|^{\frac{\gamma+2}{\gamma+1}}.$$

Then <sup>nak60</sup> (3.4) can be written in the form

$$\begin{aligned} k_2 \int_t^{t+1} \int_{\Omega} |u_t(x, s)|^{\gamma+2} dx ds \leq \|u(t)\|_E^2 + \int_t^{t+1} \int_{\Omega} \frac{k_2}{2} |u_t(x, s)|^{\gamma+2} dx ds \\ + \int_t^{t+1} \int_{\Omega} \frac{\gamma+2}{\gamma+1} \left(\frac{2}{k_2\gamma+2}\right) |f(x, s)|^{\frac{\gamma+2}{\gamma+1}} dx ds, \end{aligned}$$

$$\int_t^{t+1} \|u_t(s)\|_{L^{\gamma+2}(\Omega)}^{\gamma+2} ds \leq \frac{2}{k_2} \sup_{t \in \mathbb{R}} (\|u(t)\|_E^2 + \frac{\gamma+2}{\gamma+1} \left(\frac{2}{k_2^{\gamma+2}}\right) \|f\|_S^{\frac{\gamma+2}{\gamma+1}}).$$

Take  $M_1^{\gamma+2} = \frac{2}{k_2} \sup_{t \in \mathbb{R}} (\|u(t)\|_E^2 + \frac{\gamma+2}{\gamma+1} (\frac{2}{k_2^{\gamma+2}}) \|f\|_S^{\frac{\gamma+2}{\gamma+1}})$ , then

$$\int_t^{t+1} \|u_t(s)\|_{L^{\gamma+2}(\Omega)}^{\gamma+2} ds \leq M_1^{\gamma+2}. \quad (3.5) \quad \boxed{\text{estU}}$$

Suppose that  $u$  and  $v$  are two arbitrary solutions of the problem  $(\text{nk1})$ - $(\text{nk3})$ . Then the function  $w$  is a solution of the equation

$$w_{tt}(x, t) - \Delta w(x, t) + (|u_t(x, t)|^\gamma u_t(x, t) - |v_t(x, t)|^\gamma v_t(x, t)) = 0. \quad (3.6) \quad \boxed{\text{nak6}}$$

Multiplying  $(\text{nk6})$  by  $w_t$ , integrating over the region  $\Omega \times [t, t+1]$  and using Green's identity we get

$$\begin{aligned} \frac{d}{dt} \int_t^{t+1} \left( \frac{1}{2} \|w_t(s)\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\nabla w(s)\|_{L^2(\Omega)}^2 \right) ds \\ + \int_t^{t+1} \int_\Omega w_t(x, s) (|u_t(x, s)|^\gamma u_t(x, s) - |v_t(x, s)|^\gamma v_t(x, s)) dx ds = 0, \end{aligned}$$

$$\frac{d}{dt} \int_t^{t+1} \|w(s)\|_E^2 ds + \int_t^{t+1} \int_\Omega w_t(x, s) (|u_t(x, s)|^\gamma u_t(x, s) - |v_t(x, s)|^\gamma v_t(x, s)) dx ds = 0.$$

By Lemma  $\text{nk2L}$   $\text{I.1.18}$ , we have

$$\frac{d}{dt} \int_t^{t+1} \|w(s)\|_E^2 ds + k_1 \int_t^{t+1} \int_\Omega |w_t(x, s)|^{\gamma+2} dx ds \leq 0.$$

Integrating the last inequality over the interval  $[t, t+1]$ , we obtain

$$k_1 \int_t^{t+1} \int_\Omega |w_t(x, s)|^{\gamma+2} dx ds \leq \|w(t)\|_E^2 - \|w(t+1)\|_E^2 := k_1(A(t))^{\gamma+2}$$

where

$$A(t) = \left( \frac{1}{k_1} \|w(t)\|_E^2 - \frac{1}{k_1} \|w(t+1)\|_E^2 \right)^{\frac{1}{\gamma+2}}. \quad (3.7) \quad \boxed{\text{nak61}}$$

Using the mean value theorem for integration, we can find two points  $t_1 \in [t, t + \frac{1}{4}]$  and  $t_2 \in [t + \frac{3}{4}, t + 1]$  such that

$$\|w_t(t_i)\|_{L^{\gamma+2}(\Omega)} \leq 2A(t) \quad , \quad i = 1, 2. \quad (3.8) \quad \boxed{\text{nak8}}$$

Now we will multiply the equation  $\boxed{\text{B.6}}$  with  $w$  and integrate over  $\Omega \times [t_1, t_2]$

$$\begin{aligned} \frac{d}{dt} \int_{t_1}^{t_2} \int_{\Omega} w(x, s) w_t(x, s) dx ds - \int_{t_1}^{t_2} \int_{\Omega} w_t^2(x, s) dx ds + \int_{t_1}^{t_2} \int_{\Omega} w(x, s) \Delta w(x, s) dx ds \\ + \int_{t_1}^{t_2} \int_{\Omega} w(x, s) (|u_t(x, s)|^\gamma u_t(x, s) - |v_t(x, s)|^\gamma v_t(x, s)) dx ds = 0, \end{aligned}$$

$$\begin{aligned} \int_{t_1}^{t_2} \|w(s)\|_{H_0^1(\Omega)}^2 ds \leq |(w(t_2), w_t(t_2))| + |(w(t_1), w_t(t_1))| \\ + \left| \int_{t_1}^{t_2} (w(s), (|u_t(s)|^\gamma u_t(s) - |v_t(s)|^\gamma v_t(s))) ds \right| + \int_{t_1}^{t_2} \|w_t(s)\|_{L^2(\Omega)}^2 ds. \end{aligned}$$

By Sobolev lemma and  $\boxed{\text{B.8}}$

$$\begin{aligned} |(w(t_1), w_t(t_1))| &\leq \int_{\Omega} |w(t_1)| |w_t(t_1)| dx \leq \left( \int_{\Omega} |w(t_1)|^{\frac{\gamma+2}{\gamma+1}} dx \right)^{\frac{\gamma+1}{\gamma+2}} \left( \int_{\Omega} |w_t(t_1)|^{\gamma+2} dx \right)^{\frac{1}{\gamma+2}} \\ &\leq \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} \|w(t_1)\|_{L^{\gamma+2}(\Omega)} \|w_t(t_1)\|_{L^{\gamma+2}(\Omega)} \leq \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} 2A(t) \|w(t_1)\|_{L^{\gamma+2}(\Omega)} \\ &\leq 2\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} A(t) S_2 \|w(t_1)\|_{H_0^1(\Omega)} \leq 2\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} A(t) S_2 \|w(t_1)\|_E \\ &\leq 2\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} S_2 A(t) \max_{s \in [t, t+1]} \|w(s)\|_E. \end{aligned}$$

Similarly,

$$|(w(t_2), w_t(t_2))| \leq 2\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} S_2 A(t) \max_{s \in [t, t+1]} \|w(s)\|_E.$$

By (3.8) we have,

$$\begin{aligned} \int_{t_1}^{t_2} \|w_t(s)\|_{L^2(\Omega)}^2 ds &= \int_{t_1}^{t_2} \int_{\Omega} w_t^2(x, s) dx ds \\ &\leq \int_{t_1}^{t_2} mes(\Omega)^{\frac{\gamma}{\gamma+2}} \left( \int_{\Omega} w_t^{2 \cdot \frac{\gamma+2}{2}}(x, s) dx \right)^{\frac{2}{\gamma+2}} ds \leq mes(\Omega)^{\frac{\gamma}{\gamma+2}} A(t)^2. \end{aligned}$$

Employing the Lemma 1.1.17 we get

$$\begin{aligned} &\left| \int_{t_1}^{t_2} \int_{\Omega} (|u_t(x, s)|^\gamma u_t(x, s) - |v_t(x, s)|^\gamma v_t(x, s)) w(x, s) dx ds \right| \\ &\leq k_0 \int_{t_1}^{t_2} \int_{\Omega} (1 + |u_t(x, s)|^\gamma + |v_t(x, s)|^\gamma) |w_t(x, s)| |w(x, s)| dx ds \\ &\leq k_0 \int_{t_1}^{t_2} mes(\Omega)^{\frac{\gamma}{\gamma+2}} \left( \int_{\Omega} |w_t(x, s)|^{\frac{\gamma+2}{2} \cdot 2} dx \right)^{\frac{2}{\gamma+2} \cdot \frac{1}{2}} \left( \int_{\Omega} |w(x, s)|^{\frac{\gamma+2}{2} \cdot 2} dx \right)^{\frac{2}{\gamma+2} \cdot \frac{1}{2}} ds \\ &+ k_0 \int_{t_1}^{t_2} \|u_t(x, s)\|_{L^{\gamma+2}}^\gamma \left( \int_{\Omega} |w_t(x, s)|^{\frac{\gamma+2}{2} \cdot 2} dx \right)^{\frac{2}{\gamma+2} \cdot \frac{1}{2}} \left( \int_{\Omega} |w(x, s)|^{\frac{\gamma+2}{2} \cdot 2} dx \right)^{\frac{2}{\gamma+2} \cdot \frac{1}{2}} ds \\ &+ k_0 \int_{t_1}^{t_2} \|v_t(x, s)\|_{L^{\gamma+2}}^\gamma \left( \int_{\Omega} |w_t(x, s)|^{\frac{\gamma+2}{2} \cdot 2} dx \right)^{\frac{2}{\gamma+2} \cdot \frac{1}{2}} \left( \int_{\Omega} |w(x, s)|^{\frac{\gamma+2}{2} \cdot 2} dx \right)^{\frac{2}{\gamma+2} \cdot \frac{1}{2}} ds \\ &= k_0 \int_{t_1}^{t_2} (mes(\Omega)^{\frac{\gamma}{\gamma+2}} + \|u_t(x, s)\|_{L^{\gamma+2}}^\gamma + \|v_t(x, s)\|_{L^{\gamma+2}}^\gamma) \|w_t(x, s)\|_{L^{\gamma+2}} \|w(x, s)\|_{L^{\gamma+2}} ds. \end{aligned}$$

Due to the estimate  $\overset{\text{estU}}{(3.5)}$  and the Sobolev lemma we obtain

$$\begin{aligned}
&\leq k_0(\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma) \int_{t_1}^{t_2} \|w_t(s)\|_{L^{\gamma+2}} \|w(s)\|_{L^{\gamma+2}} ds \\
&\leq k_0(\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma) \int_{t_1}^{t_2} \|w_t(s)\|_{L^{\gamma+2}} S_{\gamma+2} \|w(s)\|_{H_0^1(\Omega)} ds \\
&\leq k_0(\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma) S_{\gamma+2} \int_{t_1}^{t_2} \|w_t(s)\|_{L^{\gamma+2}} \|w(s)\|_E ds \\
&\leq k_0(\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma) S_{\gamma+2} \max_{s \in [t, t+1]} \|w(s)\|_E \int_{t_1}^{t_2} \|w_t(s)\|_{L^{\gamma+2}} ds \\
&\leq k_0(\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma) S_{\gamma+2} \max_{s \in [t, t+1]} \|w(s)\|_E A(t).
\end{aligned}$$

Therefore,

$$\int_{t_1}^{t_2} \|w(s)\|_{H_0^1(\Omega)} \leq C_0 A(t)^2 + C_1 A(t) \max_{s \in [t, t+1]} \|w(s)\|_E \quad (3.9) \quad \boxed{\text{nak33}}$$

$$C_0 = \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}}$$

$$C_1 = 4S_2 \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + k_0 S_{\gamma+2} (\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma).$$

We use  $\overset{\text{nak61}}{(3.7)}$  and get

$$\begin{aligned}
\int_{t_1}^{t_2} \|w(s)\|_E dt &= \int_{t_1}^{t_2} \|w(s)\|_{H_0^1(\Omega)}^2 ds + \int_{t_1}^{t_2} \int_{\Omega} |w_t(s)|^2 ds \\
&\leq \int_{t_1}^{t_2} \|w(s)\|_{H_0^1(\Omega)}^2 ds + \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} \int_{t_1}^{t_2} \left( \int_{\Omega} |w_t(s)|^{\gamma+2} \right)^{\frac{2}{\gamma+2}} \\
&\leq \int_{t_1}^{t_2} \|w(s)\|_{H_0^1(\Omega)}^2 ds + \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} A(t)^2. \quad (3.10) \quad \boxed{\text{nak34}}
\end{aligned}$$

By  $\overset{\text{nak33}}{(3.9)}$  and  $\overset{\text{nak34}}{(3.10)}$ , we get

$$\begin{aligned}
\int_{t_1}^{t_2} \|w(s)\|_E &\leq \text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} A(t)^2 + C_0 A(t)^2 + C_1 A(t) \max_{s \in [t, t+1]} \|w(s)\|_E \\
&= 2C_0 A(t)^2 + C_1 A(t) \max_{s \in [t, t+1]} \|w(s)\|_E.
\end{aligned}$$



Using the mean value theorem for integration there exists  $t^* \in [t_1, t_2]$  such that

$$\|w(t^*)\|_E \leq 4C_0A(t)^2 + 2C_1A(t) \max_{s \in [t, t+1]} \|w(s)\|_E$$

When we multiply [\(3.6\)](#) by  $w_t$  and integrate  $\Omega \times [t^*, s]$  (or  $\Omega \times [t^*, s]$ ), we have

$$\pm \|w(s)\|_E \mp \|w(t^*)\|_E + \int_{t_1}^{t_2} (|u_t(s)|^\gamma u_t(s) - |v_t(s)|^\gamma v_t(s), w_t(s)) ds = 0.$$

Then  $\forall s \in [t, t+1]$  we obtain

$$\|w(s)\|_E^2 \leq \|w(t^*)\|_E^2 + \int_t^{t+1} (|u_t(s)|^\gamma u_t(s) - |v_t(s)|^\gamma v_t(s), w_t(s)) ds.$$

Thus,

$$\begin{aligned} \max_{s \in [t, t+1]} \|w(s)\|_E^2 &\leq \|w(t^*)\|_E^2 \\ &+ \int_t^{t+1} (|u_t(s)|^\gamma u_t(s) - |v_t(s)|^\gamma v_t(s), w_t(s)) ds. \end{aligned} \quad (3.11) \quad \boxed{\text{nak36}}$$

Similar to what we did above we obtain

$$\begin{aligned} \int_t^{t+1} (|u_t(s)|^\gamma u_t(s) - |v_t(s)|^\gamma v_t(s), w_t(s)) ds &\leq k_0 \int_t^{t+1} (1 + |u_t(s)|^\gamma + |v_t(s)|^\gamma) |w_t(s)| |w_t(s)| ds \\ &\leq k_0 (\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma) A(t)^2. \end{aligned}$$

Then, [\(3.11\)](#) can be written in the form

$$\begin{aligned} \max_{s \in [t, t+1]} \|w(s)\|_E^2 &\leq \|w(t^*)\|_E^2 + C_2 A(t)^2 \\ &\leq 4C_0A(t)^2 + 2C_1A(t) \max_{s \in [t, t+1]} \|w(s)\|_E + C_2A(t)^2. \end{aligned}$$

such that  $C_2 = k_0(\text{mes}(\Omega)^{\frac{\gamma}{\gamma+2}} + M_1^\gamma + M_2^\gamma)$ .

By Young's inequality, we have

$$2C_1A(t) \max_{s \in [t, t+1]} \|w(s)\|_E \leq 2C_1^2A(t)^2 + \frac{1}{2} \max_{s \in [t, t+1]} \|w(s)\|_E^2 \quad (3.12) \quad \boxed{\text{nak37}}$$

We use [\(3.12\)](#) to obtain

$$\max_{s \in [t, t+1]} \|w(s)\|_E^2 \leq (8C_0 + 2C_2 + 4C_1^2)A(t)^2.$$

Then,

$$\max_{s \in [t, t+1]} (\|w(s)\|_E^2)^{\frac{\gamma+2}{2}} \leq (8C_0 + 2C_2 + 4C_1^2)^{\frac{\gamma+2}{2}} \frac{1}{k_1} (\|w(t)\|_E^2 - \|w(t+1)\|_E^2).$$

Let's take  $K = \frac{k_1}{(8C_0 + 2C_2 + 4C_1^2)^{\frac{\gamma+2}{2}}}$  then,

$$K(\|w(s)\|_E^2)^{\frac{\gamma+2}{2}} \leq (\|w(t)\|_E^2 - \|w(t+1)\|_E^2) \quad , \forall t.$$

If we take  $\Phi = \|w(t)\|_E^2$  and  $0 < \frac{\gamma+2}{2}$  by Nakao lemma 1, we get

$$\|w(t)\|_E \leq \left(\frac{\gamma}{2}K(t-1) + M^{-\gamma}\right)^{-\frac{1}{\gamma}} \text{ for } t \geq 1.$$

And if  $0 = \frac{\gamma+2}{2}$  then by Nakao lemma 2 we get the result

$$\|w(t)\|_E \leq Me^{-K't} \quad \text{for } t \geq 1.$$

□

## Chapter 4

# An Equation of Nonlinear Vibration of Elastic String

In this chapter we study the asymptotic behavior of solutions to the initial boundary value problem for the equation of the nonlinear vibration of elastic string. We will show that the energy norm of the solution of this problem decays with an exponential rate as  $t \rightarrow +\infty$ . We also prove that the solution of the problem continuously depends on data.

The problem we consider in this chapter is:

$$u_{tt}(x, t) - (1 + \|\nabla u(t)\|^2)\Delta u(t)(x, t) + 2\gamma u_t(x, t) = f(x, t), \quad x \in \Omega, t > 0, \quad (4.1) \quad \boxed{\text{ns1}}$$

$$u(x, t) = 0, \quad x \in \partial\Omega, t > 0, \quad (4.2) \quad \boxed{\text{ns2}}$$

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad x \in \Omega. \quad (4.3) \quad \boxed{\text{ns3}}$$

We assume that the data  $u_0$ ,  $u_1$  and  $f$  are so smooth, then the problem has a classical solution. We also assume that:

$$\int_0^\infty \|f(\tau)\|^2 d\tau < \infty, \quad (4.4) \quad \boxed{\text{ns9}}$$

$$\delta_0(t) = \left( \int_t^{t+1} \|f(\tau)\|^2 d\tau \right)^{\frac{1}{2}} \leq C_0 e^{-\theta_0 t} \quad (4.5) \quad \boxed{\text{ns10}}$$

for some constants  $C_0, \theta_0 > 0$ ,

$$\left( \int_0^\infty \|\Delta f(\tau)\|^2 d\tau \right)^{\frac{1}{2}} < \infty, \quad (4.6) \quad \boxed{\text{ns11}}$$

$$\delta_{00}(t) = \left( \int_t^{t+1} \|\Delta f(\tau)\|^2 d\tau \right)^{\frac{1}{2}} \leq C_{00} e^{-\theta_{00} t} \quad (4.7) \quad \boxed{\text{ns12}}$$

for some constants  $C_{00}, \theta_{00} > 0$ .

## 4.1 Asymptotic behavior of Solutions of the Equation Nonlinear Vibration of an Elastic String

nis5L **Lemma 4.1.1.** <sup>nis2</sup> ~~(7)~~ Suppose that the assumption <sup>ns9</sup> (4.4) is satisfied. Then for solution of the problem <sup>ns1</sup> (4.1) - <sup>ns3</sup> (4.3) the following estimates hold true

$$\|u_t(t)\|, \|\nabla u(t)\|, \int_0^\infty \|u_t(\tau)\|^2 d\tau \leq C_1, \quad \forall t \geq 0,$$

where  $C_1$  is constant independent of  $t$ .

*Proof.* Let us multiply the equation <sup>ns1</sup> (4.1) by  $2u_t$  in  $L^2(\Omega)$

$$2(u_{tt}(t), u_t(t)) - 2(1 + \|\nabla u(t)\|^2)(u_t(t), \Delta u(t)) + 4\gamma \|u_t(t)\|^2 = 2(f(t), u_t(t)).$$

Using Young's inequality we get

$$\begin{aligned} \frac{d}{dt} \left( \|u_t(t)\|^2 + \|\nabla u(t)\|^2 + \frac{1}{2} \|\nabla u(t)\|^4 \right) + 4\gamma \|u_t(t)\|^2 \\ \leq \frac{1}{2\gamma} \|f(t)\|^2 + 2\gamma \|u_t(t)\|^2 \end{aligned} \quad (4.8) \quad \boxed{\text{nis14}}$$

Denoting  $E(t) = \|u_t(t)\|^2 + \|\nabla u(t)\|^2 + \frac{1}{2} \|\nabla u(t)\|^4$ , and integrating <sup>nis14</sup> (4.8) over the interval  $[0, t]$  we obtain

$$E(t) + 2\gamma \int_0^t \|u_t(\tau)\|^2 d\tau \leq \frac{1}{2\gamma} \int_0^t \|f(\tau)\|^2 d\tau + E(0).$$

Thus the inequality implies

$$\|u_t(t)\|, \|\nabla u(t)\| \text{ and } \int_0^\infty \|u_t(\tau)\|^2 d\tau \leq C_1, \quad \forall t \geq 0.$$

□

nis6L **Lemma 4.1.2.** <sup>nis2</sup> ~~(7)~~ Suppose that the assumptions <sup>ns9</sup> (4.4) and <sup>ns10</sup> (4.5) are satisfied. Then the solution of problem <sup>ns1</sup> (4.1) - <sup>ns3</sup> (4.3) satisfies

$$\|u_t(t)\|, \|\nabla u(t)\| \leq C_2 \exp(-\theta'_0 t) \text{ for some } \theta'_0 > 0, \quad \forall t \geq 0. \quad (4.9) \quad \boxed{\text{nis85}}$$

*Proof.* Integrating  $\overset{\text{nis14}}{(4.8)}$  over the interval  $(t, t+1)$  we get

$$2\gamma \int_t^{t+1} \|u_t(\tau)\|^2 d\tau \leq \frac{1}{2\gamma} \int_t^{t+1} \|f(\tau)\|^2 d\tau + E(t) - E(t+1) \equiv B(t)^2 \quad (4.10) \quad \boxed{\text{nis15}}$$

There exist two points  $t_1 \in [t, t + \frac{1}{4}]$  and  $t_2 \in [t + \frac{3}{4}, t + 1]$  such that

$$2\gamma \|u_t(t_i)\|^2 \leq 4B(t)^2 \quad i = 1, 2 \quad (4.11) \quad \boxed{\text{nis16}}$$

Multiplying  $\overset{\text{ns1}}{(4.1)}$  by  $u$  in  $L^2(\Omega)$ , integrating the obtained inequality over  $[t_1, t_2]$  and using Cauchy-Schwartz inequality, we get

$$\begin{aligned} \int_{t_1}^{t_2} (1 + \|\nabla u(\tau)\|^2) \|\nabla u(\tau)\|^2 d\tau &= \int_{t_1}^{t_2} (-(u_{tt}(\tau), u(\tau)) - 2\gamma(u_t(\tau), u(\tau)) + (f(\tau), u(\tau))) d\tau \\ &\leq |(u_t(t_1), u(t_1))| + |(u_t(t_2), u(t_2))| \\ &\quad + \int_{t_1}^{t_2} \|u_t(\tau)\|^2 d\tau + \int_{t_1}^{t_2} (2\gamma \|u_t(\tau)\| + \|f(\tau)\|) \|u(\tau)\| d\tau \end{aligned}$$

Due to inequalities  $\overset{\text{nis15}}{(4.10)}$  and  $\overset{\text{nis16}}{(4.11)}$ , we obtain

$$\begin{aligned} &\leq 2 \left(\frac{2}{\gamma}\right)^{\frac{1}{2}} B(t) \max_{\tau \in [t, t+1]} \|u(\tau)\| + \frac{B(t)^2}{2\gamma} + (2(2\gamma)^{\frac{1}{2}} B(t) + \delta_0(t)) \max_{\tau \in [t, t+1]} \|u(\tau)\| \\ &\leq 2 \left(\frac{2}{\gamma}\right)^{\frac{1}{2}} + (2\gamma)^{\frac{1}{2}} B(t) + \delta_0(t) \max_{\tau \in [t, t+1]} \|u(\tau)\| + \frac{B(t)^2}{2\gamma} \quad (4.12) \quad \boxed{\text{nis17}} \end{aligned}$$

Adding  $\overset{\text{nis15}}{(4.10)}$  and  $\overset{\text{nis17}}{(4.12)}$  we obtain,

$$\int_{t_1}^{t_2} E(\tau) d\tau \leq C \{B(t) + \delta_0(t)\} \max_{\tau \in [t, t+1]} E(\tau)^{\frac{1}{2}} + B(t)^2.$$

So there exists  $t^* \in [t_1, t_2]$  such that

$$E(t^*) \leq 2C \{B(t) + \delta_0(t)\} \max_{\tau \in [t, t+1]} E(\tau)^{\frac{1}{2}} + B(t)^2.$$

If we integrate  $\overset{\text{nis14}}{(4.8)}$  from  $t^*$  to  $s$  with  $s \in [t, t+1]$ , we get

$$\begin{aligned} E(s) &\leq E(t^*) - 2\gamma \int_{t^*}^s \|u_t(\tau)\|^2 d\tau + \frac{1}{2\gamma} \int_{t^*}^s \|f(\tau)\|^2 d\tau \\ &\leq C \left\{ (B(t) + \delta_0(t)) \max_{\tau \in [t, t+1]} E(\tau)^{\frac{1}{2}} + B(t)^2 + \delta_0(t)^2 \right\}. \end{aligned}$$

Using Young's inequality, we obtain

$$\begin{aligned} \max_{\tau \in [t, t+1]} E(\tau) &\leq C_\epsilon B(t)^2 + \epsilon \max_{\tau \in [t, t+1]} E(\tau) + C_\epsilon \delta_0(t)^2 \\ &+ \epsilon \max_{\tau \in [t, t+1]} E(\tau) + B(t)^2 + \delta_0(t)^2. \end{aligned} \quad (4.13) \quad \boxed{\text{nis60}}$$

Then  $\boxed{\text{nis60}}$  implies

$$\max_{\tau \in [t, t+1]} E(\tau) \leq C(B(t)^2 + \delta_0(t)^2) \leq C(E(t) - E(t+1)) + c\delta_0(t)^2.$$

By Lemma  $\boxed{\text{nis4L}}$  (I.1.27)

$$E(\tau) \leq C_2 \exp(-\theta'_0 t) \text{ for some } \theta'_0, C_2 > 0.$$

Hence we get

$$\|u_t(t)\|, \|\nabla u(t)\| \leq C_2 \exp(-\theta'_0 t) \text{ for some } \theta'_0, C_2 > 0.$$

□

$\boxed{\text{nis7L}}$  **Lemma 4.1.3.**  $\boxed{\text{nis2}}$  If we assume  $\boxed{\text{ns11}}$  (4.6) in addition to the assumptions of Lemma 4.1.2 for the problem  $\boxed{\text{ns1}}$  (4.1) -  $\boxed{\text{ns3}}$  (4.3), we have

$$\|\Delta u_t(t)\|, \|\Delta \nabla u(t)\|, \int_0^\infty \|\Delta u_t(\tau)\|^2 d\tau \leq C_3$$

*Proof.* Multiplying  $\boxed{\text{ns1}}$  (4.1) by  $2\Delta^2 u_t$  in  $L^2(\Omega)$  we get

$$\begin{aligned} &\frac{d}{dt} (\|\Delta u_t(t)\|^2) + \frac{d}{dt} (\|\Delta \nabla u(t)\|^2) + \frac{d}{dt} (\|\nabla u(t)\|^2 \|\Delta \nabla u(t)\|^2) \\ &- 2(\nabla u(t), \nabla u_t(t)) (\|\Delta \nabla u(t)\|^2) + 4\gamma \|\Delta u_t(t)\|^2 = 2(\Delta f(t), \Delta u_t(t)). \end{aligned} \quad (4.14) \quad \boxed{\text{nis61}}$$

Due to the inequality  $\boxed{\text{inter1}}$  (I.5) and Poincare-Friedrich's inequality we have

$$\|\nabla u_t(t)\| \leq C \|\Delta u_t(t)\| \quad (4.15) \quad \boxed{\text{nis19}}$$

By using Young's inequality, Cauchy- Schwartz inequality and  $\boxed{\text{nis19}}$  (4.15) the equation  $\boxed{\text{nis61}}$  (4.14) can be written as

$$\begin{aligned} \frac{d}{dt} E_0(t) + 2\gamma \|\Delta u_t(t)\|^2 &\leq \frac{1}{2\gamma} \|\Delta f(t)\|^2 + 2(\nabla u(t), \nabla u_t(t)) (\|\Delta \nabla u(t)\|^2) \\ &\leq \frac{1}{2\gamma} \|\Delta f(t)\|^2 + C \|\nabla u(t)\| \cdot \|\nabla u_t(t)\| \cdot E_0(t) \\ &\leq \frac{1}{2\gamma} \|\Delta f(t)\|^2 + C \|\nabla u(t)\| E_0(t)^{\frac{1}{2}} \cdot E_0(t) \end{aligned} \quad (4.16) \quad \boxed{\text{nis18}}$$

where

$$E_0(t) = \|\Delta u_t(t)\|^2 + \|\Delta \nabla u(t)\|^2 + \|\nabla u(t)\|^2 \|\Delta \nabla u(t)\|^2.$$

Integrating [\(4.16\)](#) and using [\(4.15\)](#), [\(4.9\)](#) we get

$$\begin{aligned} E_0(t) + 2\gamma \int_0^t \|\Delta u_t(\tau)\|^2 d\tau &\leq E_0(0) + \frac{1}{2\gamma} \int_0^\infty \|\Delta f(\tau)\|^2 d\tau \\ &\quad + C \int_0^t \exp(-\theta'_1 \tau) (E_0(\tau))^{\frac{1}{2}} E_0(\tau) d\tau. \end{aligned}$$

By the Bihari Lemma [1.1.21](#) we have

$$E_0(t) \leq G^{-1} \left( C \int_0^t \exp(-\theta'_1 \tau) d\tau \right) < \infty,$$

where  $G(\eta) = \int_k^\eta \frac{ds}{s^2}$ , and  $k = E_0(0) + \frac{1}{2\gamma} \int_0^\infty \|\Delta f(\tau)\|^2 d\tau < \infty$ .

Thus  $E_0(t) \leq C_3$ . Hence,

$$\|\Delta u_t(t)\|, \|\Delta \nabla u(t)\|, \int_0^\infty \|\Delta u_t(t)\|^2 dt \leq C_3.$$

□

[nis8L](#) **Theorem 4.1.4.** [Nis2](#) Suppose that [\(4.4\)](#), [\(4.5\)](#), [\(4.6\)](#) and [\(4.7\)](#) are satisfied. Then the following estimates for the solutions of [\(4.1\)](#) - [\(4.3\)](#) holds true:

$$\|\Delta u_t(t)\|, \|\Delta \nabla u(t)\| \leq C_{00} e^{-\theta_{00} t} \quad \text{for some } C_{00}, \theta_{00} > 0. \quad (4.17) \quad \text{nis100}$$

*Proof.* Integrating [\(4.16\)](#) over the interval  $(t, t+1)$  and using [\(4.15\)](#) we get,

$$\begin{aligned} 2\gamma \int_t^{t+1} \|\Delta u_t(\tau)\|^2 d\tau &\leq C \int_t^{t+1} \exp(-\theta'_1 \tau) (E_0(\tau))^{\frac{1}{2}} E_0(\tau) d\tau \\ &\quad + \frac{1}{2\gamma} \int_t^{t+1} \|\Delta f(\tau)\|^2 d\tau + E_0(t) - E_0(t+1). \quad (4.18) \quad \text{nis63} \end{aligned}$$

By Lemma [4.1.3](#), relations [\(4.7\)](#) and [\(4.18\)](#) we have

$$2\gamma \int_t^{t+1} \|\Delta u_t(\tau)\|^2 d\tau \leq C\delta_{00}^2(t) + C \exp(-\theta'_1\tau) + E_0(t) - E_0(t+1) \equiv B_0(t)^2 \quad (4.19) \quad \text{nis20}$$

Therefore, there exist two points  $t_1 \in [t, t + \frac{1}{4}]$  and  $t_2 \in [t + \frac{3}{4}, t + 1]$  such that

$$2\gamma \|\Delta u_t(t_i)\|^2 \leq 4B_0(t)^2, \quad i = 1, 2 \quad (4.20) \quad \text{nis21}$$

Multiplying [\(4.1\)](#) by  $\Delta^2 u(x, t)$  in  $L^2(\Omega)$  we get

$$2(1 + \|\nabla u(t)\|^2) \|\Delta \nabla u(t)\|^2 = -2 \frac{d}{dt} (\Delta u(t), \Delta u_t(t)) + 2\|\Delta u_t(t)\|^2 - 4\gamma (\Delta u(t), \Delta u_t(t)) + 2(\Delta u(t), \Delta f(t)) \quad (4.21) \quad \text{nis22}$$

Integrating [\(4.21\)](#) over  $(t_1, t_2)$  and using Cauchy- Schwartz inequality, we get

$$\begin{aligned} & \int_{t_1}^{t_2} 2(1 + \|\nabla u(\tau)\|^2) \|\Delta \nabla u(\tau)\|^2 d\tau \leq 2|\Delta u(t_1), \Delta u_t(t_1)| + 2|\Delta u(t_2), \Delta u_t(t_2)| \\ & + 2 \int_{t_1}^{t_2} \|\Delta u_t(\tau)\|^2 d\tau + 4\gamma \int_{t_1}^{t_2} \|\Delta u(\tau)\| \|\Delta u_t(\tau)\| d\tau + 2 \int_{t_1}^{t_2} \|\Delta u(\tau)\| \|\Delta f(\tau)\| d\tau. \end{aligned}$$

By [\(4.19\)](#) and [\(4.20\)](#), we obtain

$$\begin{aligned} & \leq CB_0(t) \max_{\tau \in [t, t+1]} \|\Delta u(\tau)\| + CB_0(t)^2 + \int_{t_1}^{t_2} (4\gamma \|\Delta u_t(\tau)\| + 2\|\Delta f(\tau)\|) \|\Delta u(\tau)\| d\tau \\ & \leq CB_0(t) \max_{\tau \in [t, t+1]} \|\Delta u(\tau)\| + CB_0(t)^2 + C(B_0(t) + \delta_{00}(t)) \max_{\tau \in [t, t+1]} \|\Delta u(\tau)\|. \end{aligned} \quad (4.22) \quad \text{nis23}$$

Adding [\(4.19\)](#) and [\(4.22\)](#) we obtain

$$E_0(t) \leq C(B_0(t) + \delta_{00}(t)) \max_{\tau \in [t, t+1]} \|\Delta u(\tau)\| + CB_0(t)^2,$$

$$\max_{\tau \in [t, t+1]} E_0(t) \leq C(B_0(t) + \delta_{00}(t)) \max_{\tau \in [t, t+1]} E_0(t)^{\frac{1}{2}} + CB_0(t)^2.$$



By Young inequality and definition of  $B_0(t)^2$ , we get

$$\max_{\tau \in [t, t+1]} E_0(t) \leq C\delta_{00}(t)^2 + C \exp(-\theta'_1 \tau) + E_0(t) - E_0(t+1)$$

By Nakao Lemma (<sup>nis4L</sup>(I.I.27)), we get

$$E_0(t) \leq C_{00} \exp(-\theta_0 t) \quad \text{for some } C_{00}, \theta_0 > 0$$

Hence,

$$\|\Delta u_t(t)\|, \|\Delta \nabla u(t)\| \leq C_{00} e^{-\theta_0 t}.$$

□

#### 4.1.1 Continuous Dependence on Data

Let us consider the problem

$$U_{tt}(x, t) - (1 + \|\nabla U(t)\|^2) \Delta U(x, t) + 2\gamma U_t(x, t) = h(x, t), \quad (4.23) \quad \boxed{\text{ns4}}$$

$$U(x, t) = 0 \quad \text{on } \partial G, \quad (4.24) \quad \boxed{\text{ns5}}$$

$$U(x, 0) = U_0, U_t(x, 0) = U_1. \quad (4.25) \quad \boxed{\text{ns6}}$$

Assume that  $f(x, t) = h(x, t) = 0$  on  $\partial\Omega$ .

<sup>nis2T</sup> **Theorem 4.1.5.** <sup>nis1</sup> <sup>[6]</sup> Let  $U$  be the solution of the problem <sup>ns4</sup>(4.23)-<sup>ns6</sup>(4.25) and  $u$  be the solution of the problem <sup>ns1</sup>(4.1)-<sup>ns3</sup>(4.3).

If  $(u_0 - U_0, u_1 - U_1, f - h)$  is sufficiently small i.e

$$\begin{aligned} & \|\Delta(u_0 - U_0)\|^2 + \|\nabla(u_1 - U_1)\|^2 \\ & + \int_0^\infty \|\nabla(f - h)(\tau)\|^2 d\tau \leq \epsilon^2 \quad \text{for some } \epsilon > 0, \quad (4.26) \quad \boxed{\text{nis25}} \end{aligned}$$

then the following estimates hold true:

$$\|\nabla(u_t(t) - U_t(t))\|^2 + \|\Delta(u(t) - U(t))\|^2 \leq C\epsilon^2, \forall t \geq 0.$$

$$\int_0^t \|\nabla(u_t(\tau) - U_t(\tau))\|^2 d\tau + \int_0^t \|\Delta(u(\tau) - U(\tau))\|^2 d\tau \leq C\epsilon^2, \forall t \geq 0.$$

*Proof.* By Lemma [4.1.1](#) we have

$$\|u_t(t)\|, \|\nabla u(t)\|, \int_0^t \|u_t(\tau)\|^2 d\tau \leq C, \forall t \geq 0$$

and

$$\|U_t(t)\|, \|\nabla U(t)\|, \int_0^t \|U_t(\tau)\|^2 d\tau \leq C, \forall t \geq 0.$$

Then  $\|v(t)\| = \|u(t) - U(t)\| \leq \|u(t)\| + \|U(t)\|$ . Then,

$$\|v_t(t)\|, \|\nabla v(t)\|, \int_0^t \|v_t(\tau)\|^2 d\tau \leq C, \forall t \geq 0,$$

where  $v(t) = u(t) - U(t)$ .

Let  $u(x, t)$  be a solution of [\(4.1\)](#) - [\(4.3\)](#). Then

$$(U_{tt}(x, t) + v_{tt}(x, t)) - (1 + \|\nabla U(t) + \nabla v(t)\|^2) \Delta(U(x, t) + v(x, t)) + 2\gamma(U_t(x, t) + v_t(x, t)) = f(x, t).$$

By [\(4.23\)](#), we obtain

$$\begin{aligned} v_{tt}(x, t) - (1 + \|\nabla v(t)\|^2 + \|\nabla U(t)\|^2) \Delta v(x, t) + 2\gamma v_t(x, t) \\ = 2(\nabla U(t), \nabla v(t)) \Delta v(x, t) \\ + (2(\nabla U(t), \nabla v(t)) + \|\nabla v(t)\|^2) \Delta U(x, t) + g(x, t), \end{aligned} \quad (4.27) \quad \boxed{\text{nis27}}$$

where  $g(x, t) = f(x, t) - h(x, t)$ ,  $v_0 = u_0 - U_0$ ,  $v_1 = u_1 - U_1$ .

Multiplying [\(4.27\)](#) by  $2v_t$  and integrate over  $\Omega$ , we get

$$\begin{aligned} \frac{d}{dt} (\|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2} \|\nabla v(t)\|^4 + \|\nabla U(t)\|^2 \|\nabla v(t)\|^2) + 4\gamma \|v_t(t)\|^2 = \\ 2(\nabla U(t), \nabla U_t(t)) \|\nabla v(t)\|^2 - 2 \frac{d}{dt} ((\nabla U(t), \nabla v(t)) \|\nabla v(t)\|^2) \\ + 2((\nabla U_t(t), \nabla v(t)) + (\nabla U(t), \nabla v_t(t))) \|\nabla v(t)\|^2 \\ + 2((\nabla U(t), \nabla v(t)) + \|\nabla v(t)\|^2) (\Delta U(t), v_t(t)) + 2(g(t), v_t(t)). \end{aligned} \quad (4.28) \quad \boxed{\text{nis28}}$$

Integrate both sides of the [\(4.28\)](#) from 0 to  $t$ , since  $-(\nabla U, \nabla v_t) =$

$(\Delta U, v_t)$  we get,

$$\begin{aligned}
& (\|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2}\|\nabla v(t)\|^4 + \|\nabla U(t)\|^2\|\nabla v(t)\|^2) \\
& - (\|v_t(0)\|^2 + \|\nabla v(0)\|^2 + \frac{1}{2}\|\nabla v(0)\|^4 + \|\nabla U(0)\|^2\|\nabla v(0)\|^2) + 4\gamma \int_0^t \|v_t(\tau)\|^2 d\tau \\
& = 2 \int_0^t (\nabla U(\tau), \nabla U_t(\tau)) \|\nabla v(\tau)\|^2 d\tau - 2(\nabla U(t), \nabla v(t)) \|\nabla v(t)\|^2 + 2(\nabla U(0), \nabla v(0)) \|\nabla v(0)\|^2 \\
& + 2 \int_0^t (\nabla U_t(\tau), \nabla v(\tau)) \|\nabla v(\tau)\|^2 d\tau + 2 \int_0^t (\nabla U(\tau), \nabla v(\tau)) (\Delta U(\tau), v_t(\tau)) d\tau + 2 \int_0^t (g(\tau), v_t(\tau)) d\tau.
\end{aligned}$$

By Hölder inequality and Young's inequality we can write our equation as:

$$\begin{aligned}
& \|\nabla v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2}\|\nabla v(t)\|^4 + \|\nabla U(t)\|^2\|\nabla v(t)\|^2 + 4\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau \\
& \leq \int_0^t (\|\nabla U(\tau)\|^2 + \|\nabla U_t(\tau)\|^2) (\|\nabla v(\tau)\|^2 + 2\|\nabla v_t(\tau)\|^2) d\tau + (\|\nabla U(t)\|^2 + \|\nabla v(t)\|^2) \|\nabla v(t)\|^2 \\
& + 2(\|\nabla U(0)\| \cdot \|\nabla v(0)\|) (\|\nabla v(0)\|^2 + 2\|v_t(0)\|^2) + (\|\nabla v(0)\|^2 + 2\|v_t(0)\|^2) \\
& + (\frac{1}{2}\|\nabla v(0)\|^2 + \|\nabla U(0)\|^2) (\|\nabla v(0)\|^2 + 2\|v_t(0)\|^2) \\
& + 2 \int_0^t \|\nabla U_t(\tau)\| \cdot \|\nabla v(\tau)\| (\|\nabla v(\tau)\|^2 + 2\|v_t(\tau)\|^2) d\tau \\
& + \int_0^t \|\nabla U(\tau)\| \cdot \|\Delta U(\tau)\| (\|\nabla v(\tau)\|^2 + 2\|v_t(\tau)\|^2) d\tau \\
& + 2\gamma \int_0^t \|v_t(\tau)\|^2 d\tau + \frac{1}{2\gamma} \int_0^t \|g(\tau)\|^2 d\tau.
\end{aligned}$$

Denote  $E_1(v(t)) = (\frac{1}{2}\|\nabla v(t)\|^2 + \|\nabla v_t(t)\|^2)$ .

By  $\frac{\text{nis100}}{(4.17)}$

$$\|v_t(t)\|, \|\nabla v(t)\|, \int_0^t \|v_t\|^2 d\tau, \|U_t(t)\|, \|\nabla U(t)\|, \int_0^t \|U_t(\tau)\|^2 d\tau \leq C \quad (4.29) \quad \boxed{\text{nis75}}$$

we have

$$\begin{aligned}
E_1(v(t)) + \frac{1}{2}(1 - \|\nabla v(t)\|^2)\|\nabla v(t)\|^2 + 2\gamma \int_0^t \|v_t(\tau)\|^2 d\tau \\
\leq C \left( E_1(v(0)) + \int_0^t \|g(\tau)\|^2 d\tau \right) + C \int_0^t c(\tau) E_1(v(\tau)) d\tau \quad (4.30) \quad \boxed{\text{nis30}}
\end{aligned}$$

where  $c(\tau) = Ce^{-\theta\tau}$ ,  $\theta > 0$ . It follows from [\(4.30\)](#) that for each  $t \leq t^* := \sup \{t : \|\nabla v(t)\|^2 \leq 1\}$  the following inequality holds true

$$E_1(v(t)) \leq C \left( E_1(v(0)) + \int_0^t \|g(\tau)\|^2 d\tau \right) + C \int_0^t c(\tau) E_1(v(\tau)) d\tau.$$

By Lemma [I.1.22](#) [ClasGron](#) Gronwall's inequality and [\(4.26\)](#) we have

$$E_1(v(t)) \leq C \left( E_1(v(0)) + \int_0^t \|g(\tau)\|^2 d\tau \right) \exp \left( C \int_0^t c(\tau) d\tau \right) \leq C\epsilon^2.$$

If we choose  $\epsilon$  sufficiently small i.e.  $C\epsilon^2 < \frac{1}{4}$ , then,

$$E_1(v(t)) = \|v_t(t)\|^2 + \frac{1}{2}\|\nabla v(t)\|^2 < \frac{1}{4} \implies \|\nabla v(t)\|^2 < \frac{1}{2}.$$

Thus for sufficiently small  $\epsilon$ 's  $t^* = \infty$ .

For sufficiently small  $\epsilon$ 's by [\(4.30\)](#) and [\(4.26\)](#) we have

$$\begin{aligned}
\|v_t(t)\|^2 + \|\nabla v(t)\|^2 + 2\gamma \int_0^t \|v_t(\tau)\|^2 d\tau \leq C \left( E_1(v(0)) + \int_0^t \|g(\tau)\|^2 d\tau \right) \\
\leq C\epsilon^2. \quad (4.31) \quad \boxed{\text{nis31}}
\end{aligned}$$

Multiply [\(4.27\)](#) by  $2\Delta v_t$  and integrate over  $\Omega$  we obtain,

$$\begin{aligned}
& -\frac{d}{dt} (\|\nabla v_t(t)\|^2 + \|\Delta v(t)\|^2 + \|\nabla v(t)\|^2 \|\Delta v(t)\|^2 + \|\nabla U(t)\|^2 \|\Delta v(t)\|^2) - 4\gamma \|\nabla v_t(t)\|^2 \\
& \quad + 2(\nabla v(t), \nabla v_t(t)) \|\Delta v(t)\|^2 + 2(\nabla U(t), \nabla U_t(t)) \|\Delta v(t)\|^2 \\
& = 2\frac{d}{dt} ((\nabla U(t), \nabla v(t)) \|\Delta v(t)\|^2) - 2((\nabla U_t(t), \nabla v(t)) + (\nabla U(t), \nabla v_t(t))) \|\Delta v(t)\|^2 \\
& \quad + 2(2(\nabla U(t), \nabla v(t)) + \|\nabla v(t)\|^2)(\Delta U(t), \Delta v_t(t)) + 2(g(t), \Delta v_t(t)). \quad (4.32) \quad \boxed{\text{nis70}}
\end{aligned}$$

Since  $\|\nabla v(t)\|^2(\Delta U(t), \Delta v_t(t)) = -(\nabla U(t), \nabla v_t(t))\|\Delta v(t)\|^2$ , we integrate (4.32) over the interval  $[0, t]$  and using Cauchy-Schwartz inequality, we get

$$\begin{aligned}
& \|\nabla v_t(t)\|^2 + \|\Delta v(t)\|^2 + \|\nabla v(t)\|^2\|\Delta v(t)\|^2 + \|\nabla U(t)\|^2\|\Delta v(t)\|^2 + 4\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau \\
& - (\|\nabla v_t(0)\|^2 + \|\Delta v(0)\|^2 + \|\nabla v(0)\|^2\|\Delta v(0)\|^2 + \|\nabla U(0)\|^2\|\Delta v(0)\|^2) \\
& \leq 2 \int_0^t \|\nabla v(\tau)\| \|\nabla v_t(\tau)\| \|\Delta v(\tau)\|^2 d\tau + 2 \int_0^t \|\nabla U(\tau)\| \|\nabla U_t(\tau)\| \|\Delta v(\tau)\|^2 d\tau \\
& - 2(\nabla U(t), \nabla v(t))\|\Delta v(t)\|^2 + 2(\nabla U(0), \nabla v(0))\|\Delta v(0)\|^2 + 2 \int_0^t \|\nabla U_t(\tau)\| \|\nabla v(\tau)\| \|\Delta v(\tau)\|^2 d\tau \\
& - 4 \int_0^t (\nabla U(\tau), \nabla v(\tau))(\Delta U(\tau), \Delta v_t(\tau)) d\tau - 2 \int_0^t (\nabla g(\tau), \nabla v_t(\tau)) d\tau.
\end{aligned}$$

By Cauchy-Schwartz inequality, Young's inequality and the identity

$$(\nabla U, \nabla v)(\Delta U, \Delta v_t) = -(\Delta U, \Delta v)(\nabla U, \nabla v_t)$$

we have

$$\begin{aligned}
& \|\nabla v_t(t)\|^2 + \|\Delta v(t)\|^2 + \|\nabla v(t)\|^2\|\Delta v(t)\|^2 + \|\nabla U(t)\|^2\|\Delta v(t)\|^2 + 4\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau \\
& \leq 2 \int_0^t \|\nabla v(\tau)\| \|\nabla v_t(\tau)\| \|\Delta v(\tau)\|^2 d\tau + 2 \int_0^t \|\nabla U(\tau)\| \|\nabla U_t(\tau)\| (\|\nabla v_t(\tau)\|^2 + \|\Delta v(\tau)\|^2) d\tau \\
& + (\|\nabla U(t)\|^2 + \|\nabla v(t)\|^2)\|\Delta v(t)\|^2 + \|\nabla v_t(0)\|^2 + \|\Delta v(0)\|^2 + \|\nabla v(0)\|^2\|\Delta v(0)\|^2 \\
& + \|\nabla U(0)\|^2\|\Delta v(0)\|^2 + (\|\nabla U(0)\|^2 + \|\nabla v(0)\|^2)\|\Delta v(0)\|^2 \\
& + 2 \int_0^t \|\nabla U_t(\tau)\| \|\nabla v(\tau)\| (\|\nabla v_t(\tau)\|^2 + \|\Delta v(\tau)\|^2) d\tau \\
& + 2 \int_0^t \|\Delta U(\tau)\| \|\nabla U(\tau)\| (\|\nabla v_t(\tau)\|^2 + \|\Delta v(\tau)\|^2) d\tau \\
& + \frac{1}{2\gamma} \int_0^t \|\nabla g(\tau)\|^2 d\tau + 2\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau. \quad (4.33) \quad \boxed{\text{nis1000}}
\end{aligned}$$

Taking  $E_2(v(t)) = \|\nabla v_t(t)\|^2 + \|\Delta v(t)\|^2$  we obtain

$$\begin{aligned}
& \|\nabla v_t(0)\|^2 + \|\Delta v(0)\|^2 + 2\|\nabla v(0)\|^2\|\Delta v(0)\|^2 + 2\|\nabla U(0)\|^2\|\Delta v(0)\|^2 \\
& \leq \|\nabla v_t(0)\|^2 + \|\Delta v(0)\|^2 + 2\|\nabla v(0)\|^2(\|\Delta v(0)\|^2 + \|\Delta v_t(0)\|^2) \\
& \quad + 2\|\nabla U(0)\|^2(\|\Delta v(0)\|^2 + \|\Delta v_t(0)\|^2) \\
& \leq CE_2(v(0)). \quad (4.34) \quad \boxed{\text{nis71}}
\end{aligned}$$

By  $(\frac{\text{nis100}}{4.17}, \frac{\text{nis71}}{4.34})$  and  $(\frac{\text{nis75}}{4.29}, \frac{\text{nis1000}}{4.33})$  can be written in the form

$$\begin{aligned}
& E_2(v(t)) + \|\nabla v(t)\|^2\|\Delta v(t)\|^2 + \|\nabla U(t)\|^2\|\Delta v(t)\|^2 + 2\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau \\
& \leq C \left( E_2(v(0)) + \int_0^\infty \|\nabla g(\tau)\|^2 d\tau \right) + \int_0^t C\epsilon \|\nabla v_t(\tau)\| \|\Delta v(\tau)\|^2 d\tau \\
& \quad + (\|\nabla U(t)\|^2 + \|\nabla v(t)\|^2)\|\Delta v(t)\|^2 + \int_0^t Cc(\tau)E_2(v(\tau))d\tau.
\end{aligned}$$

Then we get,

$$\begin{aligned}
& E_2(v(t)) + 2\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau - \int_0^t C\epsilon \|\nabla v_t(\tau)\| \|\Delta v(\tau)\|^2 d\tau \\
& \leq C \left( E_2(v(0)) + \int_0^\infty \|\nabla g(\tau)\|^2 d\tau \right) + \int_0^t Cc(\tau)E_2(v(\tau))d\tau. \quad (4.35) \quad \boxed{\text{nis32}}
\end{aligned}$$

Next, multiplying  $(\frac{\text{nis27}}{4.27})$  by  $\Delta v$  and integrate over  $\Omega$  we obtain

$$\begin{aligned}
& -\frac{d}{dt}(v_t(t), \Delta v(t)) - \|\nabla v_t(t)\|^2 + (1 + \|\nabla v(t)\|^2 + \|\nabla U(t)\|^2)\|\Delta v(t)\|^2 + \gamma \frac{d}{dt}\|\nabla v(t)\|^2 \\
& = -2(\nabla U(t), \nabla v(t))\|\Delta v\|^2 - (g(t), \Delta v(t)) - 2(\nabla U(t), \nabla v(t)) + \|\nabla v(t)\|^2(\Delta U(t), \Delta v(t)) \\
& \leq 2\|\nabla U(t)\| \|\nabla v(t)\| (\|\Delta v(t)\|^2 + \|\nabla v_t(t)\|^2) + 2\|U(t)\| \|\Delta U(t)\| (\|\Delta v(t)\|^2 + \|\nabla v_t(t)\|^2) \\
& \quad + (\|\nabla U(t)\|^2 + \|\nabla v(t)\|^2)\|\Delta v(t)\|^2 + \frac{1}{2}\|g(t)\|^2 + \frac{1}{2}\|\Delta v(t)\|^2. \quad (4.36) \quad \boxed{\text{nis76}}
\end{aligned}$$

If we integrate  $\overset{\text{nis76}}{(4.36)}$  from 0 to  $t$ , we have

$$\begin{aligned}
\int_0^t \frac{1}{2} \|\Delta v(\tau)\|^2 d\tau &\leq -(v_t(0), \Delta v(0)) + (v_t(t), \Delta v(t)) + \frac{\gamma}{2} \|\nabla v(t)\|^2 + \gamma \|\nabla v(0)\|^2 \\
&\quad + \int_0^t Cc(\tau) E_2(v(\tau)) d\tau + \frac{1}{2} \int_0^t \|g(\tau)\|^2 d\tau + \int_0^t \|\nabla v_t(\tau)\|^2 d\tau \\
&\leq \frac{1}{2} \|\nabla v_t(0)\|^2 + \left(\frac{1}{2} + \gamma\right) \|\nabla v(0)\|^2 + \frac{\gamma}{2} \|\nabla v(t)\|^2 + \frac{1}{2\gamma} \|\nabla v_t(t)\|^2 \\
&\quad + \int_0^t Cc(\tau) E_2(v(\tau)) d\tau + \frac{1}{2} \int_0^t \|g(\tau)\|^2 d\tau + \int_0^t \|\nabla v_t(\tau)\|^2 d\tau.
\end{aligned}$$

Then we have

$$\begin{aligned}
\int_0^t \frac{1}{2} \|\Delta v(\tau)\|^2 d\tau &\leq C \left( E_2 v(0) + \int_0^\infty \|g(t)\|^2 dt \right) + \frac{\gamma}{2} \|\nabla v(t)\|^2 + \frac{1}{2\gamma} \|\nabla v_t(t)\|^2 \\
&\quad + \int_0^t Cc(\tau) E_2(v(\tau)) d\tau + \int_0^t \|\nabla v_t(\tau)\|^2 d\tau. \quad (4.37) \quad \boxed{\text{nis33}}
\end{aligned}$$

If we multiply  $\overset{\text{nis33}}{(4.37)}$  by  $\gamma$  and add right hand side  $\frac{1}{2} \|\Delta v(t)\|^2$  we have

$$\begin{aligned}
\frac{\gamma}{2} \int_0^t \|\Delta v(\tau)\|^2 d\tau &\leq C\gamma \left( E_2 v(0) + \int_0^\infty \|g(\tau)\|^2 d\tau \right) + \frac{\gamma^2}{2} \|\nabla v(t)\|^2 + \gamma \int_0^t Cc(\tau) E_2(v(\tau)) d\tau \\
&\quad + \gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau + \frac{1}{2} E_2(v(t)). \quad (4.38) \quad \boxed{\text{nis34}}
\end{aligned}$$

Adding  $\overset{\text{nis34}}{(4.38)}$  to  $\overset{\text{nis32}}{(4.35)}$ , using Poincare inequality to function  $g(t)$  and  $\overset{\text{nis31}}{(4.31)}$ , we get

$$\begin{aligned}
\frac{1}{2} E_2(v(t)) + \gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau + \int_0^t \left( \frac{\gamma}{2} - C\epsilon \|\nabla v_t(\tau)\| \right) \|\Delta v(\tau)\|^2 d\tau \\
\leq C \left( E_2 v(0) + \epsilon^2 + \int_0^\infty \|\nabla g(t)\|^2 dt \right) + \int_0^t Cc(\tau) E_2(v(\tau)) d\tau. \quad (4.39) \quad \boxed{\text{nis400}}
\end{aligned}$$

If  $t^{**} = \sup \{t : \|\nabla v_t(t)\| \leq \frac{\gamma}{2C\epsilon}\} > 0$ , for  $t < t^{**}$  (4.39) is true. So we get

$$E_2(v(t)) \leq C \left( E_2v(0) + \epsilon^2 + \int_0^\infty \|\nabla g(t)\|^2 dt \right) + \int_0^t Cc(\tau)E_2(v(\tau))d\tau.$$

By (4.26) and lemma (4.1.22) Gronwall's inequality

$$E_2(v(t)) \leq C \left( E_2v(0) + \epsilon^2 + \int_0^\infty \|\nabla g(t)\|^2 d\tau \right) \exp \int_0^t Cc(\tau)d\tau \leq C\epsilon^2 \quad (4.40) \quad \text{nis36}$$

If we choose  $\epsilon$  sufficiently small, then  $E_2(v(t)) = \|\nabla v_t(t)\|^2 + \|\Delta v(t)\|^2 \leq \frac{\gamma}{4C\epsilon}$ , then  $\|\nabla v_t(t)\| < \sqrt{C\epsilon} < \frac{\gamma}{4C\epsilon}$ , which means for sufficiently small  $\epsilon$ ,  $t^{**} = \infty$ .

By (4.40) and (4.26) the inequality (4.35) can be written as:

$$\begin{aligned} E_2(v(t)) + 2\gamma \int_0^t \|\nabla v_t(\tau)\|^2 d\tau \\ \leq C \left( E_2(v(0)) + C\epsilon^2 + C \int_0^\infty \|\nabla g(t)\|^2 dt + \int_0^t Cc(\tau)E_2(v(\tau)) \right) d\tau \\ \leq C\epsilon^2 \quad (4.41) \quad \text{nis37} \end{aligned}$$

By (4.38) and (4.41) we have

$$\int_0^t \|\Delta v(\tau)\|^2 d\tau \leq C\epsilon^2 \quad (4.42) \quad \text{nis38}$$

Hence we get the results:

$$\|\nabla v_t(t)\|^2 + \|\Delta v(t)\|^2 \leq C\epsilon^2$$

$$\int_0^t \|\nabla v_t(\tau)\|^2 d\tau + \int_0^t \|\Delta v(\tau)\|^2 d\tau \leq C\epsilon^2$$

□

<sup>nis1</sup> [6] Let  $U$  be the solution of the problem (4.23) - (4.25) and  $u$  be the solution of the problem (4.1) - (4.3). If  $(u_0 - U_0, u_1 - U_1, f - h)$  is sufficiently



small in the sense  $\stackrel{\text{nis25}}{(4.26)}$  and

$$\left( \int_t^{t+1} \|\nabla(f-h)(\tau)\|^2 d\tau \right)^{\frac{1}{2}} \leq C_2 e^{-\theta't} \text{ for } \theta' > 0 \quad (4.43) \quad \boxed{\text{nis39}}$$

then we have:

$$\|\nabla(u(t) - U(t))\|^2 + \|u_t(t) - U_t(t)\|^2 \leq C_3 e^{-\theta''t}.$$

$\boxed{\text{nis3T}}$  **Theorem 4.1.6.** *Proof.* Integrating  $\stackrel{\text{nis28}}{(4.28)}$  over the interval  $[t, t+1]$  we get

$$\begin{aligned} & \|v_t(t+1)\|^2 + \|\nabla v(t+1)\|^2 + \frac{1}{2} \|\nabla v(t+1)\|^4 + \|\nabla U(t+1)\|^2 \|\nabla v(t+1)\|^2 \\ & - (\|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2} \|\nabla v(t)\|^4 + \|\nabla U(t)\|^2 \|\nabla v(t)\|^2) \\ & + 4\gamma \int_t^{t+1} \|v_t(\tau)\|^2 d\tau = 2 \int_t^{t+1} (\nabla U(\tau), \nabla U_t(\tau)) \|\nabla v(\tau)\|^2 d\tau \\ & - 2(\nabla U(t+1), \nabla v(t+1)) \|\nabla v(t+1)\|^2 + 2(\nabla U(t), \nabla v(t)) \|\nabla v(t)\|^2 \\ & + 2 \int_t^{t+1} (\nabla U_t(\tau), \nabla v(\tau)) \|\nabla v(\tau)\|^2 d\tau + 2 \int_t^{t+1} (\nabla U(\tau), \nabla v(\tau)) (\Delta U(\tau), v_t(\tau)) d\tau + 2 \int_t^{t+1} (g(\tau), v_t(\tau)) d\tau. \end{aligned}$$

By the Hölder inequality and Young's inequality

$$\begin{aligned}
& \|v_t(t+1)\|^2 + \|\nabla v(t+1)\|^2 + \frac{1}{2}\|\nabla v(t+1)\|^4 \\
& \quad - (\|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2}\|\nabla v(t)\|^4) + 4\gamma \int_t^{t+1} \|v_t(\tau)\|^2 d\tau \\
& \leq \|\nabla U(t)(\tau)\|^2 \|\nabla v(t)(\tau)\|^2 + \int_t^{t+1} (\|\nabla U(\tau)\|^2 + \|\nabla U_t(\tau)\|^2) \|\nabla v(\tau)\|^2 d\tau \\
& \quad + 2\|\nabla U(t+1)\| \|\nabla v(t+1)\| \|\nabla v(t+1)\|^2 \\
& \quad + 2\|\nabla U(t)\| \|\nabla v(t)\| \|\nabla v(t)\|^2 + 2 \int_t^{t+1} (\|\nabla U_t(\tau)\| \|\nabla v(\tau)\|) \|\nabla v(\tau)\|^2 d\tau \\
& + 2 \int_t^{t+1} \|\nabla U(\tau)\| \|\Delta U(\tau)\| \|\nabla v(\tau)\| \|v_t(\tau)\| d\tau + \frac{1}{2\gamma} \int_t^{t+1} \|g(\tau)\|^2 d\tau + 2\gamma \int_t^{t+1} \|v_t(\tau)\|^2 d\tau \\
& \leq \|\nabla U(t)\|^2 \|\nabla v(t)\|^2 + \int_t^{t+1} (\|\nabla U(\tau)\|^2 + \|\nabla U_t(\tau)\|^2) \|\nabla v(\tau)\|^2 d\tau \\
& \quad + 2\|\nabla U(t+1)\| \|\nabla v(t+1)\| \|\nabla v(t+1)\|^2 \\
& \quad + 2\|\nabla U(t)\| \|\nabla v(t)\| \|\nabla v(t)\|^2 + 2 \int_t^{t+1} (\|\nabla U_t(\tau)\| \|\nabla v(\tau)\|) \|\nabla v(\tau)\|^2 d\tau \\
& + \int_t^{t+1} \|\nabla U(\tau)\| \|\Delta U(\tau)\| (\|\nabla v(\tau)\|^2 + \|v_t(\tau)\|^2) d\tau + \frac{1}{2\gamma} \int_t^{t+1} \|g(\tau)\|^2 d\tau + 2\gamma \int_t^{t+1} \|v_t(\tau)\|^2 d\tau.
\end{aligned}$$

By [\(4.31\)](#) and [\(4.40\)](#)

$$\begin{aligned}
& \|v_t(t+1)\|^2 + \|\nabla v(t+1)\|^2 + \frac{1}{2}\|\nabla v(t+1)\|^4 - (\|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2}\|\nabla v(t)\|^4) \\
& \quad + 2\gamma \int_t^{t+1} \|v_t(\tau)\|^2 d\tau \leq Ce^{-\theta_1 t}.
\end{aligned}$$

Let's denote  $E_3(v(t)) = \|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \frac{1}{2}\|\nabla v(t)\|^4$ , then

$$2\gamma \int_t^{t+1} \|v_t(\tau)\|^2 d\tau \leq Ce^{-\theta_1 t} + E_3(t) - E_3(t+1) = B(t)^2 \quad (4.44) \quad \boxed{\text{nis40}}$$

where  $\theta_1 = \min(\theta', \theta)$ .

By [\(4.43\)](#) and Poincare-Friedrich's inequality, we obtain

$$\left( \int_t^{t+1} \|g(\tau)\|^2 d\tau \right)^{\frac{1}{2}} \leq \frac{1}{\lambda_1} \left( \int_t^{t+1} \|\nabla g(\tau)\|^2 d\tau \right)^{\frac{1}{2}} \leq C e^{-\theta t}.$$

By [\(4.44\)](#) there exist two points  $t_1 \in [t, t + \frac{1}{4}]$ ,  $t_2 \in [t + \frac{3}{4}, t + 1]$  such that

$$2\gamma \|v_t(t_i)\|^2 \leq 4B(t)^2 \quad i = 1, 2. \quad \text{nis41}$$

By Theorem [4.1.4](#)

$$\int_t^{t+1} \|\Delta U(\tau)\| d\tau \leq C e^{-\theta t}. \quad \text{nis42}$$

Multiply [\(4.27\)](#) by  $v$ , then we get

$$\begin{aligned} & \frac{d}{dt} (v(t), v_t(t)) - \|v_t(t)\|^2 + \|\nabla v(t)\|^2 + \|\nabla v(t)\|^4 + \|\nabla v(t)\|^2 \|\nabla U(t)\|^2 + 2\gamma (v(t), v_t(t)) \\ &= -2(\nabla U(t), \nabla v(t)) \|\nabla v(t)\|^2 + 2((\nabla U(t), \nabla v(t)) + \|\nabla v(t)\|^2)(\Delta U(t), v(t)) + (g(t), v(t)). \end{aligned}$$

Then,

$$\begin{aligned} & \|\nabla v(t)\|^2 + \|\nabla v(t)\|^4 + \|\nabla v(t)\|^2 \|\nabla U(t)\|^2 \\ & \leq -\frac{d}{dt} (v(t), v_t(t)) + \|v_t(t)\|^2 + 4\|\nabla U(t)\| \cdot \|\nabla v(t)\| \cdot \|\nabla v(t)\|^2 \\ & \quad + (g(t), v(t)) - 2\gamma (v(t), v_t(t)) + 2\|\nabla v(t)\|^2 \|\nabla U(t)\|^2. \end{aligned}$$

By [\(4.31\)](#) and  $\|\nabla v_t(t)\|^2, \|\Delta v(t)\|^2 \leq C\epsilon^2$

$$\begin{aligned} \|\nabla v(t)\|^2 + \|\nabla v(t)\|^4 & \leq -\frac{d}{dt} (v(t), v_t(t)) + \|v_t(t)\|^2 \\ & \quad + C e^{-\theta t} + 2\gamma \|v(t)\| \|v_t(t)\| + \|v(t)\| \|g(t)\|. \quad \text{nis43} \end{aligned}$$

If we integrate [\(4.47\)](#) over  $[t_1, t_2]$ , we have

$$\begin{aligned} \int_{t_1}^{t_2} (1 + \|\nabla v(\tau)\|^2) \|\nabla v(\tau)\|^2 d\tau & \leq |v(t_1), v_t(t_1)| + |v(t_2), v_t(t_2)| + \int_{t_1}^{t_2} \|v_t(\tau)\|^2 d\tau \\ & \quad + \int_{t_1}^{t_2} (2\gamma \|v_t(\tau)\| + \|g(\tau)\|) d\tau + C e^{-\theta t}. \end{aligned}$$

Since  $Ce^{-\theta t} < CB(t)^2$  we have

$$\begin{aligned}
& \int_{t_1}^{t_2} (1 + \|\nabla v(\tau)\|^2) \|\nabla v(\tau)\|^2 d\tau \\
& \leq \max_{t \leq \tau \leq t+1} \|v(\tau)\| (|v_t(t_1)| + |v_t(t_2)|) + CB(t)^2 \\
& + \max_{t \leq \tau \leq t+1} \|v(\tau)\| \int_{t_1}^{t_2} (2\gamma \|v_t(\tau)\| + \|g(\tau)\|) d\tau + CB(t)^2 \\
& \leq CB(t) \max_{t \leq \tau \leq t+1} \|v(\tau)\| + CB(t)^2. \quad (4.48) \quad \boxed{\text{nis44}}
\end{aligned}$$

Adding  $\boxed{\text{nis40}}$  and  $\boxed{\text{nis43}}$ , we get

$$\int_{t_1}^{t_2} E_3(\tau) d\tau \leq CB(t) \max_{t \leq \tau \leq t+1} E_3(\tau)^{\frac{1}{2}} + CB(t)^2.$$

Thus there exists a point  $t^* \in [t_1, t_2]$  such that

$$E_3(t^*) \leq 2C \left\{ B(t) \max_{t \leq \tau \leq t+1} E_3(\tau)^{\frac{1}{2}} + B(t)^2 \right\}. \quad (4.49) \quad \boxed{\text{nis45}}$$

We integrate  $\boxed{\text{nis28}}$  from  $t^*$  to  $s \in [t, t+1]$ :

$$\begin{aligned}
& \|v_t(s)\|^2 + \|\nabla v(s)\|^2 + \frac{1}{2} \|\nabla v(s)\|^4 \\
& = \|v_t(t^*)\|^2 + \|\nabla v(t^*)\|^2 + \frac{1}{2} \|\nabla v(t^*)\|^4 + \int_{t^*}^s -\frac{d}{dt} (\|\nabla U(\tau)\|^2 \|\nabla v(\tau)\|^2) d\tau \\
& - \int_{t^*}^s 4\gamma \|v_t(\tau)\| d\tau + 2 \int_{t^*}^s (\nabla U(\tau), \nabla U_t(\tau)) \|\nabla v(\tau)\|^2 d\tau - 2 \int_{t^*}^s \frac{d}{dt} (\nabla U(\tau), \nabla v(\tau)) \|\nabla v(\tau)\|^2 d\tau \\
& + 2 \int_{t^*}^s (\nabla U_t(\tau), \nabla v(\tau)) \|\nabla v(\tau)\|^2 d\tau - 2 \int_{t^*}^s (\nabla U(\tau), \nabla v(\tau)) (\Delta U(\tau), v_t(\tau)) d\tau + 2 \int_{t^*}^s (g(\tau), v_t(\tau)) d\tau.
\end{aligned}$$

Then

$$\begin{aligned}
E_3(s) &\leq E_3(t^*) + \|\nabla U(t^*)\|^2 + \|\nabla v(t^*)\|^2 + \|\nabla U(t^*)\|^2 + \|\nabla v(t^*)\|^2 \\
&+ \int_{t^*}^s (\|\nabla U(\tau)\|^2 + \|\nabla U_t(\tau)\|^2) \|\nabla v(\tau)\|^2 d\tau + 2\|\nabla U(s)\| \|\nabla v(s)\| \|\nabla v(s)\|^2 \\
&+ 2\|\nabla U(t^*)\| \|\nabla v(t^*)\| \|\nabla v(t^*)\|^2 + 2 \int_{t^*}^s \|\nabla U_t(\tau)\| \|\nabla v(\tau)\| \|\nabla v(\tau)\|^2 d\tau \\
&+ 2 \int_{t^*}^s \|\nabla U(\tau)\| \|\Delta U(\tau)\| \|\nabla v(\tau)\| \|v_t(\tau)\| d\tau + 2 \int_{t^*}^s \|g(\tau)\| \|v_t(\tau)\| d\tau.
\end{aligned}$$

By [\ref{nis45}](#), [Theorem 4.1.4](#), [\ref{nis2T}](#) and

$$\|v_t(t)\|, \|\nabla v(t)\|, \int_0^t \|v_t(\tau)\|^2 d\tau, \|U_t(t)\|, \|\nabla U(t)\|, \int_0^t \|U_t(\tau)\|^2 d\tau \leq C,$$

we get

$$E_3(s) \leq E_3(t^*) + Ce^{-\theta t} + Ce^{-\theta' t}.$$

$$E_3(s) \leq C \left\{ B(t) \max_{t \leq \tau \leq t+1} E_3(\tau)^{\frac{1}{2}} + B(t)^2 \right\} \quad \forall s \in [t, t+1],$$

$$\max_{t \leq \tau \leq t+1} E_3(\tau) \leq C \left\{ B(t) \max_{t \leq \tau \leq t+1} E_3(\tau)^{\frac{1}{2}} + B(t)^2 \right\},$$

$$\max_{t \leq \tau \leq t+1} E_3(\tau) \leq \frac{1}{2} \max_{t \leq \tau \leq t+1} E_3(\tau) + CB(t)^2.$$

Then,

$$\max_{t \leq \tau \leq t+1} E_3(\tau) \leq 2CB(t)^2 \leq C(E_3(t) - E_3(t+1) + Ce^{-\theta_1 t}).$$

By [Lemma 1.1.27](#) we get,

$$\|\nabla v(t)\|^2 + \|v_t(t)\|^2 \leq C_3 e^{-\theta'' t} \text{ for some } \theta'' > 0.$$

□

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